

Proof of Divergence Thm

Same as Green's Thm, we'll prove only the case of special domain

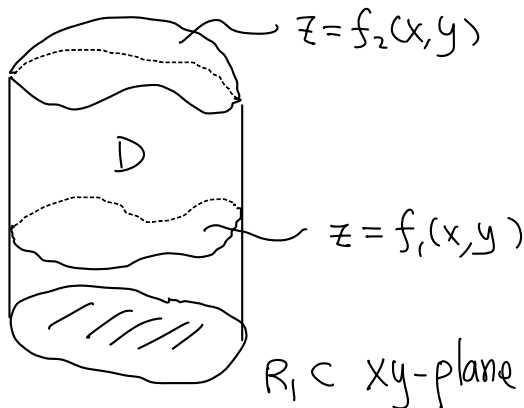
D which is of type I, II, & III:

$$D = \{(x, y, z) \in \mathbb{R}^3 : (x, y) \in R_1, f_1(x, y) \leq z \leq f_2(x, y)\} \text{ (type I)}$$

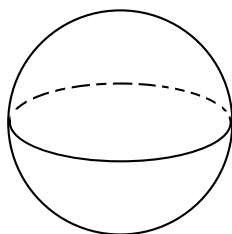
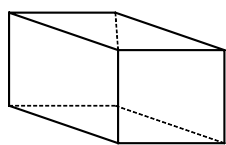
$$= \{(x, y, z) \in \mathbb{R}^3 : (y, z) \in R_2, g_1(y, z) \leq x \leq g_2(y, z)\} \text{ (type II)}$$

$$= \{(x, y, z) \in \mathbb{R}^3 : (x, z) \in R_3, h_1(x, z) \leq y \leq h_2(x, z)\} \text{ (type III)}$$

eg. type I domain



eg: special domains:



And also as in the proof of Green's Thm for

$$\vec{F} = M\hat{i} + N\hat{j} + L\hat{k}$$

we'll prove 3 equalities in the following which

combine to give the Divergence Thm:

$$\left\{ \begin{array}{l} \iint_S M \hat{i} \cdot \hat{n} \, d\sigma = \iiint_D \frac{\partial M}{\partial x} \, dV \quad (\text{by type II}) \\ \iint_S N \hat{j} \cdot \hat{n} \, d\sigma = \iiint_D \frac{\partial N}{\partial y} \, dV \quad (\text{by type II}) \\ \iint_S L \hat{k} \cdot \hat{n} \, d\sigma = \iiint_D \frac{\partial L}{\partial z} \, dV \quad (\text{by type I}) \end{array} \right.$$

The proofs are similar, we'll prove only the last one

$$\iint_S L \hat{k} \cdot \hat{n} \, d\sigma = \iiint_D \frac{\partial L}{\partial z} \, dV$$

By Fubini's Thm

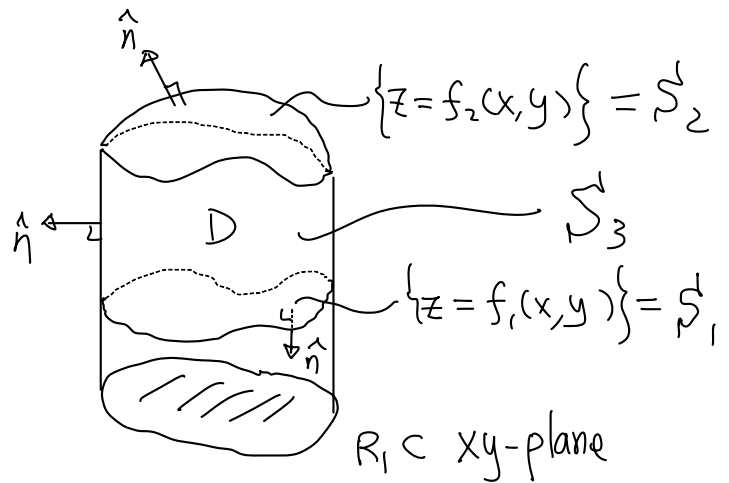
$$\begin{aligned} \text{RHS} &= \iiint_D \frac{\partial L}{\partial z} \, dV = \iint_{R_1} \left[\int_{f_1(x,y)}^{f_2(x,y)} \frac{\partial L}{\partial z} \, dz \right] dx dy \quad (\text{because Type I}) \\ &= \iint_{R_1} [L(x,y, f_2(x,y)) - L(x,y, f_1(x,y))] \, dx dy \end{aligned}$$

For the LHS, we note that

by definition of type I domain,

the boundary surface $S (= \partial D)$

of D can be written as

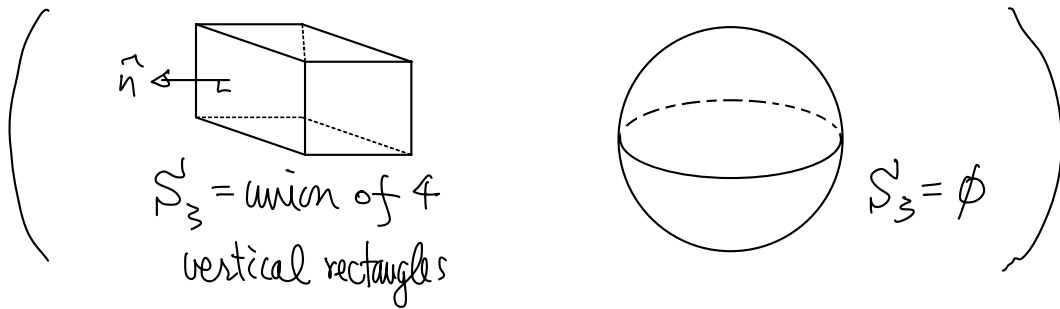


$$S = S_1 \cup S_2 \cup S_3$$

where $S_1 = \text{graph of } f_1 = \{(x, y, f_1(x, y))\} = \{z = f_1(x, y)\}$

$S_2 = \text{graph of } f_2 = \{(x, y, f_2(x, y))\} = \{z = f_2(x, y)\}$

$S_3 = \text{vertical surface (which could be empty)}$
between S_1 & S_2



$$\text{Hence LHS} = \iint_S L \hat{k} \cdot \hat{n} \, d\sigma = \iint_{S_1} L \hat{k} \cdot \hat{n} \, d\sigma + \iint_{S_2} L \hat{k} \cdot \hat{n} \, d\sigma + \iint_{S_3} L \hat{k} \cdot \hat{n} \, d\sigma$$

(Since \hat{n} of a vertical surface is horizontal, hence $\hat{k} \cdot \hat{n} = 0$)

Now on the upper surface $S_2 = \{z = f_2(x, y)\}$

the outward normal \hat{n} is upward (ie $\hat{n} \cdot \hat{k} \geq 0$)

Note that the parametrization

$$(x, y) \mapsto \vec{r}(x, y) = (x, y, f_2(x, y))$$

$$\text{has } \begin{cases} \vec{r}_x = \hat{i} + \frac{\partial f_2}{\partial x} \hat{k} \\ \vec{r}_y = \hat{j} + \frac{\partial f_2}{\partial y} \hat{k} \end{cases}$$

and $\vec{r}_x \times \vec{r}_y = -\frac{\partial f_2}{\partial x} \hat{i} - \frac{\partial f_2}{\partial y} \hat{j} + \hat{k}$ $\vec{r}_x \times \vec{r}_y$ is upward

$\Rightarrow \hat{n} = \frac{\vec{r}_x \times \vec{r}_y}{\|\vec{r}_x \times \vec{r}_y\|}$ is the upward unit normal

and $\hat{k} \cdot \hat{n} = \frac{1}{\|\vec{r}_x \times \vec{r}_y\|}$.

Therefore $\iint_{S_2} L \hat{k} \cdot \hat{n} d\sigma = \iint_{R_1} L(x, y, f_2(x, y)) \frac{1}{\|\vec{r}_x \times \vec{r}_y\|} \|\vec{r}_x \times \vec{r}_y\| dx dy$
 $= \iint_{R_1} L(x, y, f_2(x, y)) dx dy$

Similarly, note that the outward normal on S_1 (lower surface) is downward (i.e. $\hat{n} \cdot \hat{k} < 0$), we have

$\hat{n} = -\frac{\vec{r}_x \times \vec{r}_y}{\|\vec{r}_x \times \vec{r}_y\|}$, where $\vec{r}(x, y) = (x, y, f_1(x, y))$
 $(= x\hat{i} + y\hat{j} + f_1(x, y)\hat{k})$

$\Rightarrow \hat{k} \cdot \hat{n} = -\frac{1}{\|\vec{r}_x \times \vec{r}_y\|}$

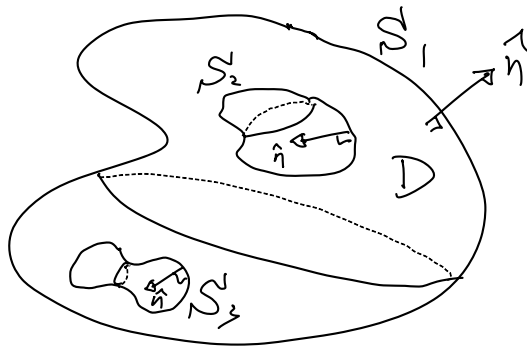
Hence $\iint_{S_1} L \hat{k} \cdot \hat{n} d\sigma = -\iint_{R_1} L(x, y, f_1(x, y)) dx dy$

$\therefore \iint_S L \hat{k} \cdot \hat{n} d\sigma = \iint_{R_1} L(x, y, f_2(x, y)) dx dy - \iint_{R_1} L(x, y, f_1(x, y)) dx dy$

$$\begin{aligned}
 &= \iint_{R_1} [L(x,y, f_z(x,y)) - L(x,y, f_i(x,y))] dx dy \\
 &= \iiint_D \frac{\partial L}{\partial z} dV.
 \end{aligned}$$

This completes the proof of the divergence thm. ~~✘~~

Note: Similar to Green's Thm, the Divergence Thm also holds for solid region with finitely many holes inside:

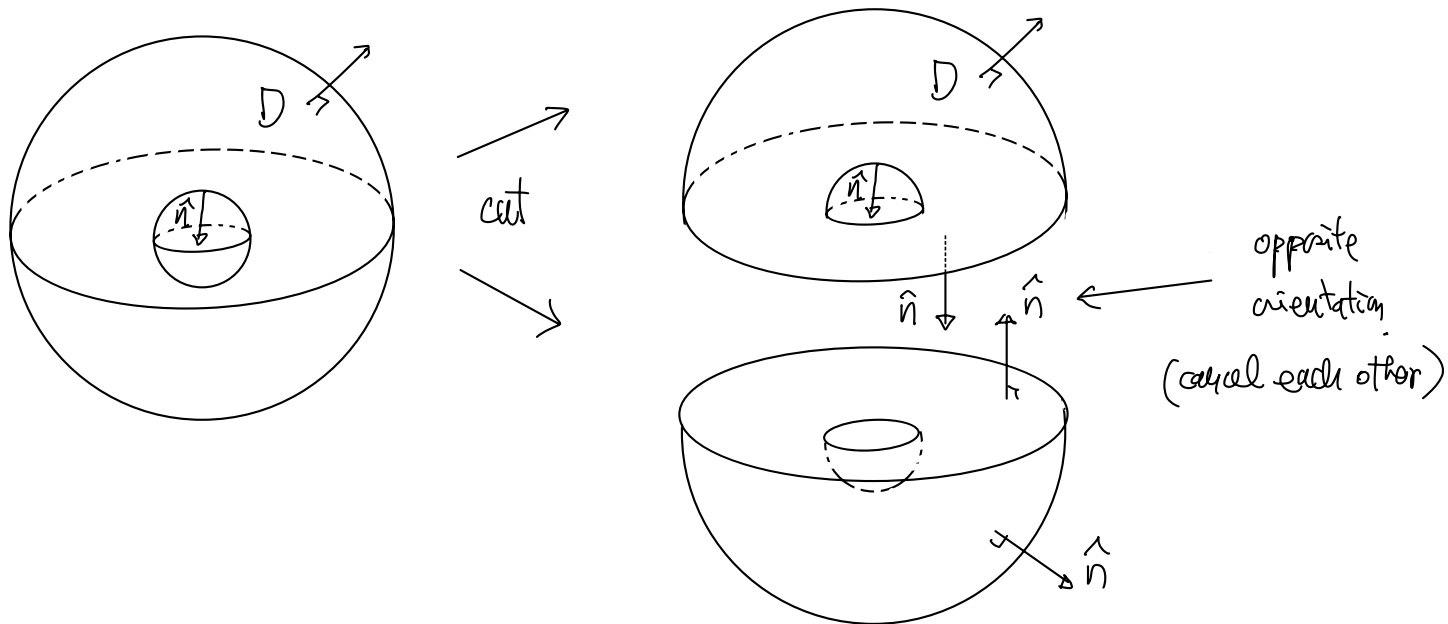


D = solid region inside S_1
but outside of S_2 and S_3

$$\iiint_D \vec{\nabla} \cdot \vec{F} dV = \sum_{i=1}^n \iint_{S_i} \vec{F} \cdot \hat{n} d\sigma$$

for \hat{n} = outward normal with respect to D .

eg (idea of proof of this kind of surface :)



Note: Physical meaning of $\text{div } \vec{F} = \vec{\nabla} \cdot \vec{F}$ in \mathbb{R}^3
= flux density (by the divergence theorem)

Unified treatment of Green's, Stokes', and Divergence Theorems

Stokes' Thm in notations of differential forms (in \mathbb{R}^3)

Making definition of differential forms

(1) A differential 1-form (or simply 1-form)

is a linear combination of the symbols dx , dy & dz :

$$\omega = \omega_1 dx + \omega_2 dy + \omega_3 dz$$

with coefficients $\omega_1, \omega_2, \omega_3$ functions on \mathbb{R}^3 .

eg: Total different of a function f on \mathbb{R}^3 :

$$df = \frac{\partial f}{\partial x} dx + \frac{\partial f}{\partial y} dy + \frac{\partial f}{\partial z} dz \quad \text{is a 1-form.}$$

(2) Wedge product: Let " \wedge " be an operation such that

$$\left\{ \begin{array}{l} dx \wedge dx = dy \wedge dy = dz \wedge dz = 0 \\ dx \wedge dy = -dy \wedge dx \\ dy \wedge dz = -dz \wedge dy \\ dz \wedge dx = -dx \wedge dz \end{array} \right.$$

and satisfies other usual rules in arithmetic.

i.e. If $\omega = \omega_1 dx + \omega_2 dy + \omega_3 dz$

$$\eta = \eta_1 dx + \eta_2 dy + \eta_3 dz$$

then we have

$$\begin{aligned}\omega \wedge \eta &= (\omega_1 dx + \omega_2 dy + \omega_3 dz) \wedge (\eta_1 dx + \eta_2 dy + \eta_3 dz) \\ &= (\omega_1 dx) \wedge (\eta_1 dx) + (\omega_2 dy) \wedge (\eta_1 dx) + (\omega_3 dz) \wedge (\eta_1 dx) \\ &\quad + (\omega_1 dx) \wedge (\eta_2 dy) + (\omega_2 dy) \wedge (\eta_2 dy) + (\omega_3 dz) \wedge (\eta_2 dy) \\ &\quad + (\omega_1 dx) \wedge (\eta_3 dz) + (\omega_2 dy) \wedge (\eta_3 dz) + (\omega_3 dz) \wedge (\eta_3 dz) \\ &= (\omega_1 \eta_1) \cancel{dx \wedge dx} + \omega_2 \eta_1 dy \wedge dx + \omega_3 \eta_1 dz \wedge dx \\ &\quad + (\omega_1 \eta_2) dx \wedge dy + \omega_2 \eta_2 \cancel{dy \wedge dy} + \omega_3 \eta_2 dz \wedge dy \\ &\quad + (\omega_1 \eta_3) dx \wedge dz + \omega_2 \eta_3 dy \wedge dz + \omega_3 \eta_3 \cancel{dz \wedge dz}\end{aligned}$$

$$\begin{aligned}\therefore \omega \wedge \eta &= (\omega_2 \eta_3 - \omega_3 \eta_2) dy \wedge dz \\ &\quad + (\omega_3 \eta_1 - \omega_1 \eta_3) dz \wedge dx \\ &\quad + (\omega_1 \eta_2 - \omega_2 \eta_1) dx \wedge dy\end{aligned}$$

- Linear combinations of $dy \wedge dz$, $dz \wedge dx$ & $dx \wedge dy$ are called differential 2-forms (on \mathbb{R}^3)

$$\zeta = \zeta_1 dy \wedge dz + \zeta_2 dz \wedge dx + \zeta_3 dx \wedge dy$$

Similarly, if ω is a 1-form and

ζ is a 2-form

then we can define $\omega \wedge \zeta$

eg: If $\omega = dx$, $\xi = dy \wedge dz$

$$\text{then } \omega \wedge \xi = dx \wedge dy \wedge dz$$

Note that we insist on the anti-commutativity of wedge product,

we have

$$\begin{aligned} dx \wedge dy \wedge dz &= - dy \wedge dx \wedge dz \\ &= dy \wedge dz \wedge dx \\ &= - dz \wedge dy \wedge dx \\ &= dz \wedge dx \wedge dy \\ &= - dx \wedge dz \wedge dy \end{aligned}$$

And $dx \wedge dx \wedge dy = \dots = 0$ whenever one of the dx, dy, dz is repeated.

Hence, as $\dim \mathbb{R}^3 = 3$, all "linear combinations" of "3-forms" are just $f dx \wedge dy \wedge dz$

which is called a differential 3-form (also called a volume form if $f > 0$)

Note: It is convenient to call smooth functions f the differential 0-form.

Summary (on \mathbb{R}^3)

$$0\text{-form} = f$$

$$1\text{-form} = \omega_1 dx + \omega_2 dy + \omega_3 dz$$

$$2\text{-form} = \xi_1 dy \wedge dz + \xi_2 dz \wedge dx + \xi_3 dx \wedge dy$$

$$3\text{-form} = g dx \wedge dy \wedge dz$$

where, f, g, ω_i, ξ_i are (smooth) functions

Note = One can certainly define k -form for any $k \geq 0$. But in \mathbb{R}^3 , k -forms are zero for $k > 3$:

$$dx^i \wedge dx \wedge dy \wedge dz = 0, \text{ where } dx^i = dx, dy, \text{ or } dz.$$