

Proof of Divergence Theorem

Same as Green's Thm, we'll prove only the case of special domain D which is of type I, II and III:

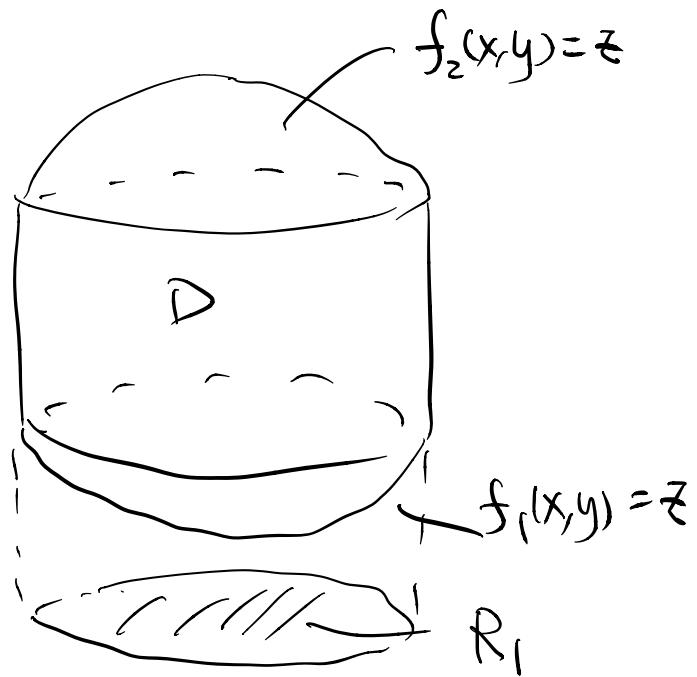
i.e.

$$D = \{(x, y, z) \in \mathbb{R}^3 : (x, y) \in R_1, f_1(x, y) \leq z \leq f_2(x, y)\} \quad (\text{type I})$$

$$= \{(x, y, z) \in \mathbb{R}^3 : (y, z) \in R_2, g_1(y, z) \leq x \leq g_2(y, z)\} \quad (\text{type II})$$

$$= \{(x, y, z) \in \mathbb{R}^3 : (z, x) \in R_3, h_1(z, x) \leq y \leq h_2(z, x)\} \quad (\text{type III})$$

e.g. of type I domain =



And also as in the proof of Green's Theorem,

$$\text{for } \vec{F} = M \hat{i} + N \hat{j} + L \hat{k},$$

we'll prove 3 equalities in the following which

combine to give the divergence thm:

$$\left\{ \begin{array}{l} \iint_S \hat{M}_i \cdot \vec{n} d\sigma = \iiint_D \frac{\partial M}{\partial x} dv \quad (\text{by type II}) \\ \iint_S \hat{N}_j \cdot \vec{n} d\sigma = \iiint_D \frac{\partial N}{\partial y} dv \quad (\text{by type III}) \\ \iint_S \hat{L}_k \cdot \vec{n} d\sigma = \iiint_D \frac{\partial L}{\partial z} dv \quad (\text{by type I}) \end{array} \right.$$

The proofs are similar, we'll prove only the last one:

$$\iint_S \hat{L}_k \cdot \vec{n} d\sigma = \iiint_D \frac{\partial L}{\partial z} dv.$$

By Fubini's Thm,

$$\text{R.H.S.} = \iiint_D \frac{\partial L}{\partial z} dv = \iint_{R_1} \left[\int_{f_1(x,y)}^{f_2(x,y)} \frac{\partial L}{\partial z} dz \right] dx dy$$

$$= \iint_{R_1} [L(x,y, f_2(x,y)) - L(x,y, f_1(x,y))] dx dy$$

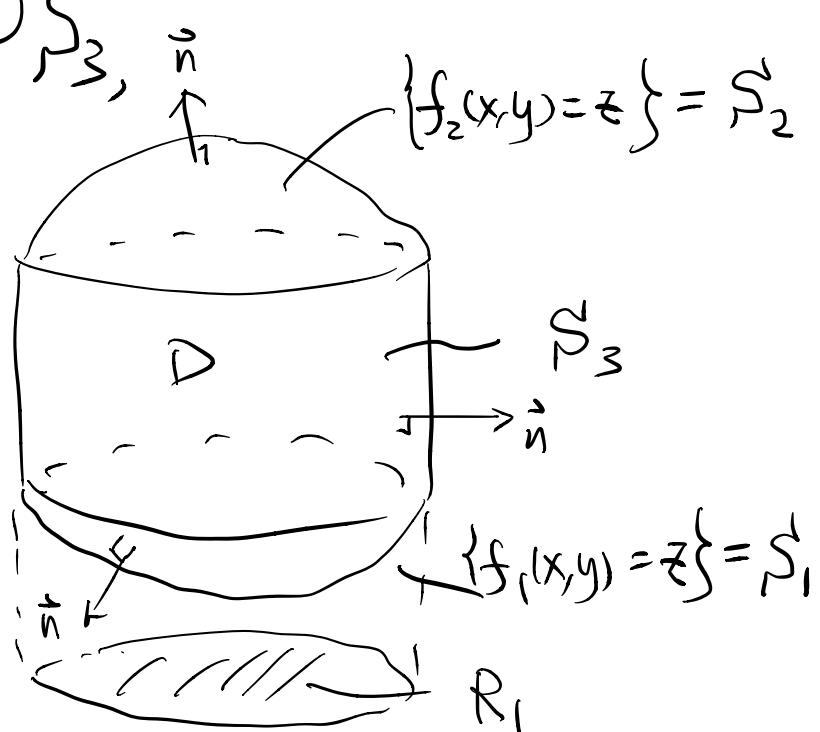
For the L.H.S., we note by definition of type I domain, the boundary S of D can be written as

$$S = S_1 \cup S_2 \cup S_3, \quad \vec{n}$$

$\{f_2(x,y) = z\} = S_2$

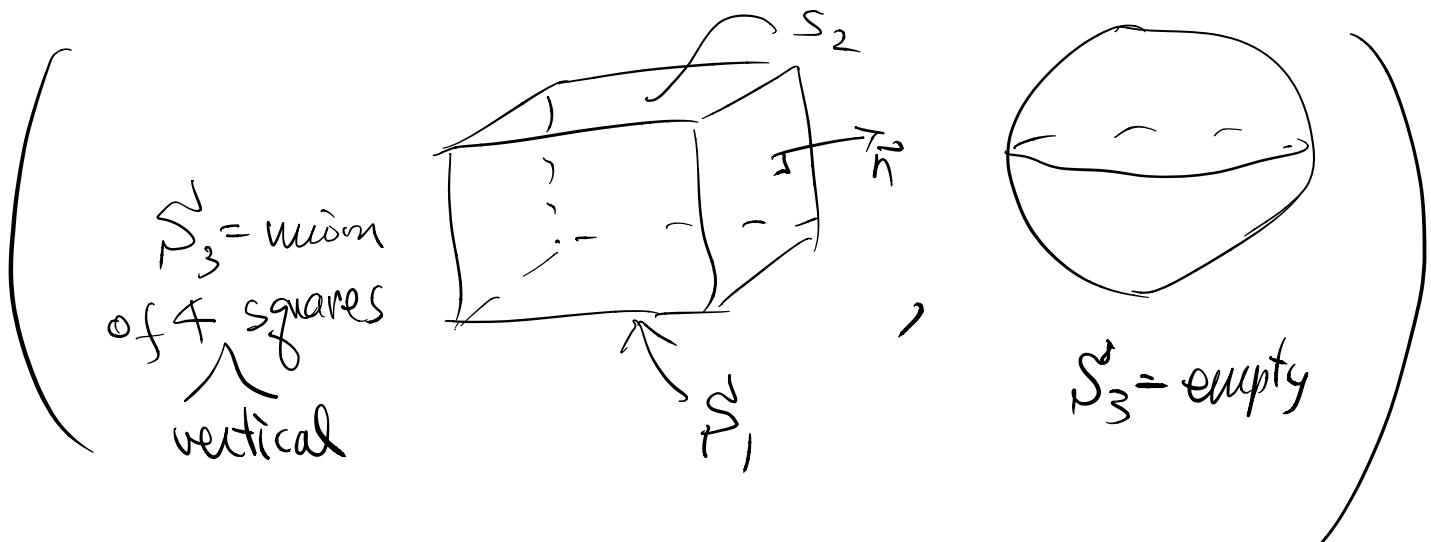
where

$$\begin{aligned} S_1 &= \text{graph of } f_1 \\ &= \{(x,y, f_1(x,y))\} \\ &= \{z = f_1(x,y)\}, \end{aligned}$$



$$\begin{aligned} S_2 &= \text{graph of } f_2 \\ &= \{(x,y, f_2(x,y))\} = \{z = f_2(x,y)\}, \text{ and} \end{aligned}$$

$S_3 = \text{a vertical surface (which could be empty!)}$



Hence

$$\text{L.H.S.} = \iint_S L \hat{k} \cdot \vec{n} d\sigma = \iint_{S_1} L \hat{k} \cdot \vec{n} d\sigma + \iint_{S_2} L \hat{k} \cdot \vec{n} d\sigma \\ + \iint_{S_3} L \hat{k} \cdot \vec{n} d\sigma$$

(Since \vec{n} of vertical surface is horizontal!)

Now on the upper surface $S_2 = \{z = f_2(x, y)\}$,
the outward normal \vec{n} is upward (in the sense that
 $\vec{n} \cdot \hat{k} > 0$). Note that the parametrization

$$(x, y) \mapsto \vec{r}(x, y) = x \hat{i} + y \hat{j} + f_2(x, y) \hat{k}$$

has
$$\begin{cases} \vec{r}_x = \hat{i} + \frac{\partial f_2}{\partial x} \hat{k} \\ \vec{r}_y = \hat{j} + \frac{\partial f_2}{\partial y} \hat{k} \end{cases}$$

$$\Rightarrow \vec{r}_x \times \vec{r}_y = \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ 1 & 0 & \frac{\partial f_2}{\partial x} \\ 0 & 1 & \frac{\partial f_2}{\partial y} \end{vmatrix} = -\frac{\partial f_2}{\partial x} \hat{i} - \frac{\partial f_2}{\partial y} \hat{j} + \hat{k}$$

$\vec{r}_x \times \vec{r}_y \stackrel{+ve}{\Rightarrow}$ is upward

$$\therefore \vec{n} = \frac{\vec{r}_x \times \vec{r}_y}{|\vec{r}_x \times \vec{r}_y|}$$

$$\text{and } \hat{k} \cdot \vec{n} = \frac{1}{|\vec{r}_x \times \vec{r}_y|}$$

$$\begin{matrix} \hat{k} \cdot \vec{n} \\ \downarrow \\ d\sigma \end{matrix}$$

Hence

$$\iint_{S_2} L \hat{k} \cdot \vec{n} d\sigma = \iint_{R_1} L(x, y, f_2(x, y)) \frac{1}{|\vec{r}_x \times \vec{r}_y|} \cdot |\vec{r}_x \times \vec{r}_y| dA$$

$$= \iint_{R_1} L(x, y, f_2(x, y)) dx dy .$$

Similarly, note that the outward normal at S_1 (lower surface) is downward (i.e. $\vec{n} \cdot \hat{k} < 0$)

$$\text{we have } \vec{n} = - \frac{\vec{r}_x \times \vec{r}_y}{|\vec{r}_x \times \vec{r}_y|}, \text{ where } \vec{r} = x\hat{i} + y\hat{j} + f_1(x, y)\hat{k}$$

$$\therefore \hat{k} \cdot \vec{n} = - \frac{1}{|\vec{r}_x \times \vec{r}_y|} . \quad (\text{check!})$$

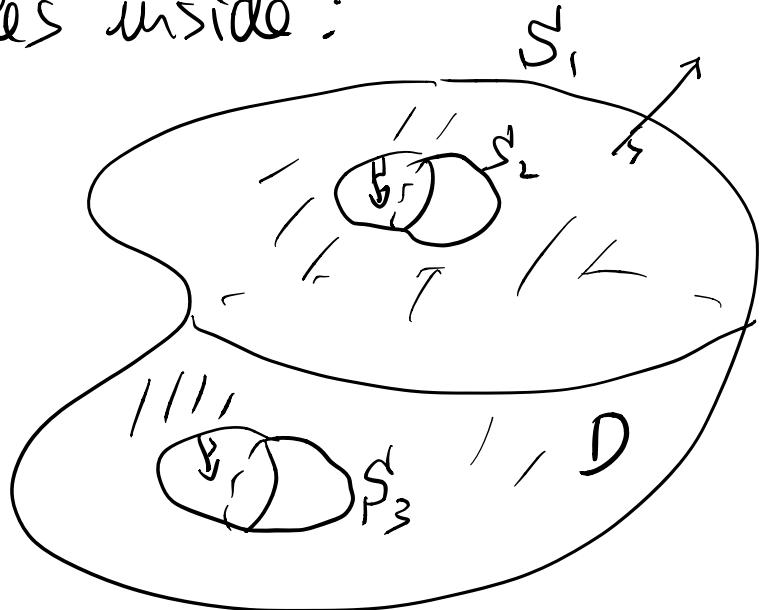
$$\text{Hence } \iint_{S_1} L \hat{k} \cdot \vec{n} d\sigma = - \iint_{R_1} L(x, y, f_1(x, y)) dx dy$$

Therefore,

$$\iint_S L \hat{k} \cdot \vec{n} d\sigma = \iint_{R_1} [L(x, y, f_2(x, y)) - L(x, y, f_1(x, y))] dx dy$$
$$= \iiint_D \frac{\partial h}{\partial z} dV . \quad \times$$

Note: Similar to Green's Theorem, the Divergence Thm is also hold for solid region with finitely many holes inside:

$$\iiint_D \vec{r} \cdot \vec{F} dV$$
$$= \sum_{i=1}^k \iint_{S_i} \vec{F} \cdot \vec{n} d\sigma$$



for \vec{n} outward normal with respect to D .

Note = Physical meaning of $\operatorname{div} \vec{F} = \vec{\nabla} \cdot \vec{F}$ in \mathbb{R}^3
= flux density (by the divergence theorem)

Unified treatment of Green's, Stokes', and Divergence Theorems :

Stokes' Theorem in notations of differential forms ($\text{in } \mathbb{R}^3$)

Working definition of differential forms:

(1) A differential 1-form (or simply 1-form) is a linear combination of the symbols dx , dy , and dz :

$$\boxed{\omega = \omega_1 dx + \omega_2 dy + \omega_3 dz}$$

with coefficients $\omega_1, \omega_2, \omega_3$ are functions on \mathbb{R}^3 .

e.g.: The total differential of a smooth function f is a differential 1-form:

$$df = \frac{\partial f}{\partial x} dx + \frac{\partial f}{\partial y} dy + \frac{\partial f}{\partial z} dz .$$

(2) Wedge product: let " \wedge " be an operation such that

$$\left\{ \begin{array}{l} dx \wedge dx = dy \wedge dy = dz \wedge dz = 0 \\ dx \wedge dy = -dy \wedge dx \\ dy \wedge dz = -dz \wedge dy \\ dz \wedge dx = -dx \wedge dz \end{array} \right.$$

and usual rules in arithmetic:

i.e. If $\omega = \omega_1 dx + \omega_2 dy + \omega_3 dz$
 $\eta = \eta_1 dx + \eta_2 dy + \eta_3 dz,$

then

$$\begin{aligned} \omega \wedge \eta &= (\omega_1 dx + \omega_2 dy + \omega_3 dz) \wedge (\eta_1 dx + \eta_2 dy + \eta_3 dz) \\ &= \cancel{\omega_1 dx \wedge \eta_1 dx} + \omega_2 dy \wedge \cancel{\eta_1 dx} + \omega_3 dz \wedge \cancel{\eta_1 dx} \\ &\quad + \omega_1 dx \wedge \cancel{\eta_2 dy} + \cancel{\omega_2 dy \wedge \eta_2 dy} + \omega_3 dz \wedge \cancel{\eta_2 dy} \\ &\quad + \omega_1 dx \wedge \cancel{\eta_3 dz} + \omega_2 dy \wedge \cancel{\eta_3 dz} + \cancel{\omega_3 dz \wedge \eta_3 dz} \\ &= (\omega_1 \eta_2 - \omega_2 \eta_1) dx \wedge dy \\ &\quad + (\omega_2 \eta_3 - \omega_3 \eta_2) dy \wedge dz \\ &\quad + (\omega_3 \eta_1 - \omega_1 \eta_3) dz \wedge dx \end{aligned}$$

$$\begin{aligned} \omega \wedge \eta &= (\omega_2 \eta_3 - \omega_3 \eta_2) dy \wedge dz \\ &\quad + (\omega_3 \eta_1 - \omega_1 \eta_3) dz \wedge dx \\ &\quad + (\omega_1 \eta_2 - \omega_2 \eta_1) dx \wedge dy \end{aligned}$$

- Linear combinations of $dy \wedge dz$, $dz \wedge dx$, & $dx \wedge dy$ are called differential 2-fams (on \mathbb{R}^3):

$$\xi = \xi_1 dy \wedge dz + \xi_2 dz \wedge dx + \xi_3 dx \wedge dy$$