

## Solutions to Assignment Two

Exercise 2.1

2. For  $\phi(x) = \log(1 + x^2)$  and  $\psi(x) = 4 + x$ , the solution is

$$\begin{aligned} u(x, t) &= \frac{1}{2} \{ \log[1 + (x + ct)^2] + \log[1 + (x - ct)^2] \} + \frac{1}{2c} \int_{x-ct}^{x+ct} 4 + s \, ds \\ &= \frac{1}{2} \{ \log[1 + (x + ct)^2] + \log[1 + (x - ct)^2] \} + 4t + xt \end{aligned}$$

5. We have

$$u(x, t) = \frac{1}{2c} \int_{x-ct}^{x+ct} \psi(s) \, ds = \frac{1}{2c} [\text{length of } (x - ct, x + ct) \cap (-a, a)].$$

Then for  $t = \frac{a}{2c}$ ,

$$u(x, t) = \frac{1}{2c} [\text{length of } (x - \frac{a}{2}, x + \frac{a}{2}) \cap (-a, a)]$$

$$= \begin{cases} 0, & x > \frac{3a}{2}; \\ \frac{3a}{2} - x, & \frac{3a}{2} \geq x > \frac{a}{2}; \\ a, & \frac{a}{2} \geq x > -\frac{a}{2}; \\ x + \frac{3a}{2}, & -\frac{a}{2} \geq x > -\frac{3a}{2}; \\ 0, & -\frac{3a}{2} \geq x. \end{cases}$$

for  $t = \frac{a}{c}$ ,

$$u(x, t) = \frac{1}{2c} [\text{length of } (x - a, x + a) \cap (-a, a)]$$

$$= \begin{cases} 0, & x > 2a; \\ 2a - x, & 2a \geq x > 0; \\ 2a + x, & 0 \geq x > -2a; \\ 0, & -2a \geq x. \end{cases}$$

for  $t = \frac{3a}{2c}$ ,

$$u(x, t) = \frac{1}{2c} [\text{length of } (x - \frac{3a}{2}, x + \frac{3a}{2}) \cap (-a, a)]$$

$$= \begin{cases} 0, & x > \frac{5a}{2}; \\ \frac{5a}{2} - x, & \frac{5a}{2} \geq x > \frac{a}{2}; \\ 2a, & \frac{a}{2} \geq x > -\frac{a}{2}; \\ x + \frac{5a}{2}, & -\frac{a}{2} \geq x > -\frac{5a}{2}; \\ 0, & -\frac{5a}{2} \geq x. \end{cases}$$

for  $t = \frac{2a}{c}$ ,

$$u(x, t) = \frac{1}{2c} [\text{length of } (x - 2a, x + 2a) \cap (-a, a)]$$

$$= \begin{cases} 0, & x > 3a; \\ 3a - x, & 3a \geq x > a; \\ 2a, & a \geq x > -a; \\ x + 3a, & -a \geq x > -3a; \\ 0, & -3a \geq x. \end{cases}$$

for  $t = \frac{5a}{c}$ ,

$$u(x, t) = \frac{1}{2c} [\text{length of } (x - 5a, x + 5a) \cap (-a, a)]$$

$$= \begin{cases} 0, & x > 6a; \\ 6a - x, & 6a \geq x > 4a; \\ 2a, & 4a \geq x > -4a; \\ x + 6a, & -4a \geq x > -6a; \\ 0, & -6a \geq x. \end{cases}$$

Omit the figures

6.

$$\max_x \{u(x, t)\} = \begin{cases} 2ct, & \frac{a}{c} > t \geq 0; \\ 2a, & t \geq \frac{a}{c}. \end{cases}$$

7. We have

$$u(x, t) = \frac{1}{2}[\phi(x + ct) + \phi(x - ct)] + \frac{1}{2c} \int_{x-ct}^{x+ct} \psi(s) ds,$$

so

$$\begin{aligned} u(-x, t) &= \frac{1}{2}[\phi(-x + ct) + \phi(-x - ct)] - \frac{1}{2c} \int_{-x-ct}^{-x+ct} \psi(-s) ds \\ &= \frac{1}{2}[-\phi(x - ct) - \phi(x + ct)] - \frac{1}{2c} \int_{x-ct}^{x+ct} \psi(s) ds \\ &= -\left\{ \frac{1}{2}[\phi(x + ct) + \phi(x - ct)] + \frac{1}{2c} \int_{x-ct}^{x+ct} \psi(s) ds \right\} = -u(x, t). \end{aligned}$$

Thus  $u(x, t)$  is also odd in  $x$  for all  $t$ .

8.(a) Change variables  $v = ru$ , then

$$\begin{aligned} v_{tt} &= ru_{tt}, \quad v_{rr} = (ru_r + u)_r = ru_{rr} + 2u_r \\ &\Rightarrow v_{tt} = rc^2(u_{rr} + \frac{2}{r}u_r) = c^2v_{rr} \end{aligned}$$

(b) Using the same skill related to the wave equation (1), we have

$$v(r, t) = f(r + ct) + g(r - ct),$$

where  $f$  and  $g$  are two arbitrary functions of a single variable. So

$$u = \frac{1}{r}f(r + ct) + \frac{1}{r}g(r - ct).$$

(c) Since  $v(r, 0) = r\phi(r)$  and  $v_t(r, 0) = r\psi(r)$  are both odd. We extend  $v$  to all of  $\mathbb{R}$  by odd reflection. That is, we set

$$\tilde{v}(r, t) := \begin{cases} v(r, t), & r > 0, t \geq 0; \\ 0, & r = 0, t \geq 0; \\ -v(-r, t), & r < 0, t \geq 0. \end{cases}$$

Hence d'Alembert's formula implies

$$\tilde{v}(r, t) = \frac{1}{2}[(r + ct)\phi(r + ct) + (r - ct)\phi(r - ct)] - \frac{1}{2c} \int_{r-ct}^{r+ct} s\psi(s) ds.$$

Therefore

$$u(r, t) = \frac{1}{r}v(r, t) = \frac{1}{2r}[(r + ct)\phi(r + ct) + (r - ct)\phi(r - ct)] - \frac{1}{2cr} \int_{r-ct}^{r+ct} s\psi(s) ds$$

for  $r > 0, t > 0$

9. Using the same skill related to the wave equation (1), let  $v = u_x + u_t$

$$\Rightarrow v_x - 4v_t = 0$$

$$\Rightarrow v(x, t) = h(x + \frac{1}{4}t)$$

$$\Rightarrow u_x + u_t = h(x + \frac{1}{4}t)$$

It is easy to check that one solution is

$$u(x, t) = f(x + \frac{1}{4}t), \text{ where } f'(s) = \frac{4h(s)}{5}.$$

So the general solution is

$$u(x, t) = f(x + \frac{1}{4}t) + g(x - t)$$

Since  $x^2 = u(x, 0) = f(x) + g(x)$  and  $e^x = u_t(x, 0) = \frac{1}{4}f'(x) - g'(x)$

$$\Rightarrow f(x) = \frac{4}{5}(x^2 + e^x + a), \quad g(x) = \frac{1}{5}(x^2 - 4e^x - 4a)$$

$$\Rightarrow u(x, t) = \frac{4}{5}[(x + \frac{t}{4})^2 + e^{x+\frac{t}{4}} + a] + \frac{1}{5}[(x - t)^2 - 4e^{x-t} - 4a]$$

$$= \frac{1}{5} \left[ \left(2x + \frac{t}{2}\right)^2 + (x-t)^2 + 4e^{x+\frac{t}{4}} - 4e^{x-t} \right]$$

Exercise 2.2

1. By the law of conservation of energy, we have

$$E = \frac{1}{2} \int_{-\infty}^{+\infty} (\rho u_t^2 + T u_x^2) dx$$

is a constant independent of  $t$ .

Since  $\phi \equiv 0$  and  $\psi \equiv 0$ , we have  $E \equiv 0$ .

The first vanishing theorem implies

$$u_t \equiv 0, u_x \equiv 0.$$

Thus  $u \equiv 0$  since  $\phi \equiv 0$ .

2. (a) Since

$$\partial e / \partial t = u_t u_{tt} + u_x u_{xt}$$

$$\partial e / \partial x = u_t u_{tx} + u_x u_{xx}$$

$$\partial p / \partial t = u_t u_{xt} + u_{tt} u_x$$

$$\partial p / \partial x = u_t u_{xx} + u_{tx} u_x$$

and

$$u_{tt} = u_{xx}$$

$$u_{xt} = u_{tx},$$

we have

$$\partial e / \partial t = \partial p / \partial x$$

$$\partial e / \partial x = \partial p / \partial t$$

(b) By (a), we have

$$e_{tt} = p_{xt} = p_{tx} = e_{xx},$$

$$p_{tt} = e_{xt} = e_{tx} = p_{xx},$$

so both  $e(x, t)$  and  $p(x, t)$  also satisfy the wave equation.

3.(a) Let  $v(x, t) = u(x - y, t)$ , we have

$$v_{tt}(x, t) = u_{tt}(x - y, t) = c^2 u_{xx}(x - y, t) = c^2 v_{xx}(x, t),$$

so it is also a solution.

(b) Let  $v(x, t) = u_x$ , we have

$$v_{tt}(x, t) = u_{xtt}(x, t) = c^2 u_{xxx}(x, t) = c^2 v_{xx}(x, t),$$

so it is also a solution.

(c) Let  $v(x, t) = u(ax, at)$ , we have

$$v_{tt}(x, t) = a^2 u_{tt}(ax, at) = a^2 c^2 u_{xx}(ax, at) = c^2 v_{xx}(x, t),$$

so it is also a solution.

5. The energy function is

$$E = \frac{1}{2} \int_{-\infty}^{+\infty} (u_t^2 + c^2 u_x^2) dx,$$

so

$$\begin{aligned} dE/dt &= \frac{1}{2} \int_{-\infty}^{+\infty} (2u_t u_{tt} + 2c^2 u_x u_{xt}) dx \\ &= \int_{-\infty}^{+\infty} (c^2 u_t u_{xx} - r u_t^2 + c^2 u_x u_{xt}) dx \\ &= \int_{-\infty}^{+\infty} (c^2 u_t u_{xx} - r u_t^2 - c^2 u_{xx} u_t) dx + c^2 u_t u_x |_{-\infty}^{+\infty} \\ &= -r \int_{-\infty}^{+\infty} (u_t^2) dx \leq 0 \end{aligned}$$

Hence the energy decreases.

More on 2.2

Part(a)

$$\begin{aligned} \frac{d}{dt} \int_0^l u^2(x, t) dx &= \int_0^l 2u u_t dx = 2k \int_0^l u u_{xx} dx \\ &= 2k u u_x |_0^l - 2k \int_0^l u_x^2 dx = 2k \left( u(l, t) u_x(l, t) - u(0, t) u_x(0, t) \right) - 2k \int_0^l u_x^2 dx \\ &= 2k \left( -\frac{u_x^2(l, t)}{a_l} - \frac{u_x^2(0, t)}{a_0} \right) - 2k \int_0^l u_x^2 dx \\ &\leq 0. \end{aligned}$$

Part(b)

Let  $u_1$  and  $u_2$  are two solutions of the above problem. Let  $w = u_1 - u_2$ , then  $w$  satisfies the same equation and the same Robin boundary condition, especially  $w(x, 0) \equiv 0$ .

By Part(a) we know that

$$\int_0^l w^2(x, t) dx \leq \int_0^l w^2(x, 0) dx = 0.$$

Thus  $w \equiv 0$ , i.e.  $u_1 \equiv u_2$ . Thus the solution of the problem is unique.

### Exercise 2.3

1.  $u(x, t) = 1 - x^2 - 2kt$ , then

$$u(0, 0) = \max_{0 \leq x \leq 1, 0 \leq t \leq T} u(x, t) = 0$$

$$u(1, T) = \min_{0 \leq x \leq 1, 0 \leq t \leq T} u(x, t) = -2kT.$$

2.(a) By the definition of Maximum,  $M(T)$  increase (i.e., nondecreasing);

(b) By the definition of minimum,  $M(T)$  decrease (i.e., nondecreasing).

3.(a) Use the strong minimum principle;

(b) Use the minimum principle or Let the maximum occur at the point  $X(t)$ , so that  $\mu(t) = u(X(t), t)$ , Differentiate  $\mu(t)$ , assuming that  $X(t)$  is differentiable, we have

$$\dot{\mu} = u_x(X(t), t)X'(t) + u_t(X(t), t) \leq 0.$$

Hence  $\mu(t)$  decrease;

(c) Omit the figure.

4.(a) Use the strong minimum principle;

(b) Since both  $u(x, t)$  and  $u(1 - x, t)$  are the solution of

$$u_t = u_{xx}, \quad 0 < x < 1, \quad 0 < t < \infty$$

$$u(0, t) = u(1, t) = 0, \quad \text{and } u(x, 0) = 4x(1 - x),$$

the uniqueness theorem implies that  $u(x, t) = u(1 - x, t)$ ;

(c) We have

$$d/dt \int_0^1 u^2 dx = \int_0^1 2uu_t dx = 2 \int_0^1 uu_{xx} dx = -2 \int_0^1 u_x^2 dx.$$

Since  $u(x, t) > 0$  for all  $t > 0$  and  $0 < x < 1$ , so  $u_x$  is not zero function.

Hence

$$d/dt \int_0^1 u^2 dx < 0.$$

Therefore  $\int_0^1 u^2 dx$  is a strictly decreasing function of  $t$ .

6. Set  $w = u - v$  and use the maximum principle.

7.(a) Set  $w = u - v$  and use the generalized Maximum Principle (which will be proved in Extra(2)(a)):

If  $u_t - u_{xx} \leq 0$  in  $R = [0, l] \times [0, T]$ , then

$$\max_R u(x, t) = \max_{x=0, x=l, t=0} u(x, t).$$

(b) Let  $u(x, t) = (1 - e^{-t}) \sin x$ , then  $u(0, t) = 0, u(\pi, t) = 0, u(x, 0) = 0, u_t - u_{xx} = \sin x$ . Hence part(a) show that  $v(x, t) \geq (1 - e^{-t}) \sin x$ .

8. We have

$$\begin{aligned} d/dt \int_0^l u^2(x, t) dx &= \int_0^l 2uu_t dx = 2k \int_0^l uu_{xx} dx = 2kuu_x|_0^l - 2k \int_0^l u_x^2 dx \\ &= 2k[u(l, t)u_x(l, t) - u(0, t)u_x(0, t)] - 2k \int_0^l u_x^2 dx \\ &= 2k\left[-\frac{u_x^2(l, t)}{a_l} - \frac{u_x^2(0, t)}{a_0}\right] - 2k \int_0^l u_x^2 dx \leq 0, \end{aligned}$$

so  $\int_0^l u^2(x, t) dx$  decrease.

Extra1. **Note:** In(1), the conclusion " $0 < u(x, t) < 1$ " has changed to " $0 < u(x, t) < e^{at}$ ". If the conclusion does not change, the condition of  $a$  must limited to  $a \leq \pi^2 k$ .

(1) Let  $v(x, t) = e^{-at}u(x, t)$ , then

$$v_t = v_{xx}$$

$$v(0, t) = v(1, t) = 0$$

$$v(x, 0) = \sin(\pi x)$$

Using the strong maximum principle, we have  $0 < v(x, t) < 1$ , therefore

$$0 < u(x, t) < e^{at}$$

for all  $t > 0$  and  $0 < x < 1$ .

(2) Let  $v(x, t) = e^{-at}u(x, t)$ , then both  $v(x, t)$  and  $v(1 - x, t)$  are solutions of the diffusion equation  $v_t = v_{xx}$  with the boundary conditions:

$$v(0, t) = v(1, t) = 0$$

$$v(x, 0) = \sin(2x).$$

Hence the uniqueness theorem implies that  $v(x, t) = v(1 - x, t)$ . Therefore

$$u(x, t) = u(1 - x, t)$$

for all  $t \geq 0$  and  $0 \leq x \leq 1$ .

Another Proof:

Let  $v(x, t) = e^{-at}u(x, t)$ , then

$$v_t = v_{xx}$$

$$v(0, t) = v(1, t) = 0$$

$$v(x, 0) = \sin(2x)$$

Using the method of separation of variables, we set  $v(x, t) = X(x)T(t)$ , then

$$X\dot{T} - k\ddot{X}T = 0$$

$$\Rightarrow \frac{\ddot{X}}{X} = -\lambda = \frac{\dot{T}}{kT}$$

$$\Rightarrow \ddot{X} + \lambda X = 0, \quad \dot{T} = -\lambda kT.$$

Hence by the boundary conditions

$$X(0) = X(1) = 0, \quad X(x)T(0) = v(x, 0) = \sin(\pi x),$$

we have

$$\begin{aligned} v(x, t) &= e^{-\pi^2 kt} \sin(\pi x) \\ \Rightarrow u(x, t) &= e^{at - \pi^2 kt} \sin(\pi x) \end{aligned}$$

Therefore  $0 < u(x, t) < e^{at}$  for all  $t > 0$  and  $0 < x < 1$ , and  $u(x, t) = u(1 - x, t)$  for all  $t \geq 0$  and  $0 \leq x \leq 1$ .

Extra2.(a) Let  $\epsilon$  be a positive constant and let

$$v(x, t) = u(x, t) + \epsilon x^2.$$

This function  $v$  satisfies

$$v_t - kv_{xx} = u_t - k(u + \epsilon x^2)_{xx} = u_t - ku_{xx} - 2\epsilon k \leq -2\epsilon k < 0,$$

which is the "diffusion inequality." Now suppose that  $v(x, t)$  attains its maximum at an interior point or on the top edge  $(x_0, t_0)$ . That is,  $0 < x_0 < 1$ ,  $0 < t_0 \leq T$ . By ordinary calculus, we know that

$$v_x(x_0, t_0) = 0, \quad v_{xx}(x_0, t_0) \leq 0$$

and

$$v_t(x_0, t_0) \geq 0$$

This contradicts the diffusion inequality. So

$$\begin{aligned} v(x, t) &\leq M + \epsilon l^2 \\ \Rightarrow u(x, t) &\leq M + \epsilon(l^2 - x^2), \\ \Rightarrow u(x, t) &\leq M \end{aligned}$$

where  $M := \max_{x=0, x=l, t=0} u(x, t)$ .

Therefore

$$\max_R u(x, t) = \max_{x=0, x=l, t=0} u(x, t).$$

(b) Consider

$$u(x, t) = v(x, t) - t \max_{-\infty < x < \infty, 0 < t < T} f(x, t),$$

then

$$\begin{aligned} u_t - k u_{xx} &= - \max_{-\infty < x < \infty, 0 < t < T} f(x, t) + f(x, t) \leq 0, \\ u(x, 0) &= 0. \end{aligned}$$

Hence part (a) or the proof of (a) implies that  $u(x, t) \leq 0$ . Therefore

$$v(x, t) \leq T \max_{-\infty < x < \infty, 0 < t < T} f(x, t).$$

#### Exercise 2.4

1. By the general formula, we have

$$\begin{aligned} u(x, t) &= \frac{1}{\sqrt{4\pi kt}} \int_{-l}^l e^{-(x-y)^2/4kt} dy \\ &= \frac{1}{\sqrt{\pi}} \int_{-l-x}^{l-x} e^{-p^2} dp \\ &= 1/2 \{ \mathfrak{Erf}[(x+l)/(2\sqrt{kt})] - \mathfrak{Erf}[(x-l)/(2\sqrt{kt})] \}. \end{aligned}$$

2. By the general formula, we have

$$\begin{aligned} u(x, t) &= \frac{1}{\sqrt{4\pi kt}} \int_0^\infty e^{-(x-y)^2/4kt} dy + \frac{1}{\sqrt{4\pi kt}} \int_{-\infty}^0 3e^{-(x-y)^2/4kt} dy \\ &= 1/2 + 1/2 \mathfrak{Erf}[x/(2\sqrt{kt})] + 3/2 - 3/2 \mathfrak{Erf}[x/(2\sqrt{kt})] \\ &= 2 - \mathfrak{Erf}[x/(2\sqrt{kt})]. \end{aligned}$$

4. By the general formula, we have

$$\begin{aligned} u(x, t) &= \frac{1}{\sqrt{4\pi kt}} \int_0^\infty e^{-(x-y)^2/4kt} e^{-y} dy \\ &= \frac{1}{\sqrt{4\pi kt}} \int_0^\infty e^{-\frac{(y+2kt-x)^2}{4kt} + kt-x} dy \\ &= e^{kt-x} \int_{\frac{2kt-x}{2\sqrt{kt}}}^\infty e^{-p^2} dp / \sqrt{\pi} = \frac{1}{2} e^{kt-x} - \frac{e^{kt-x}}{2} \mathfrak{Erf}\left(\frac{2kt-x}{2\sqrt{kt}}\right). \end{aligned}$$

6. Since

$$\begin{aligned}\int_0^\infty e^{-x^2} dx \cdot \int_0^\infty e^{-y^2} dy &= \int_0^\infty \int_0^\infty e^{-(x^2+y^2)} dx dy \\ &= \frac{\pi}{2} \int_0^\infty e^{-r^2} r dr = \frac{\pi}{4} (-e^{-r^2}) \Big|_0^\infty = \frac{\pi}{4},\end{aligned}$$

Hence

$$\int_0^\infty e^{-x^2} dx = \frac{\sqrt{\pi}}{2}.$$

7. We have

$$\begin{aligned}\int_{-\infty}^\infty e^{-p^2} dp &= \int_0^\infty e^{-p^2} dp + \int_{-\infty}^0 e^{-p^2} dp \\ &= 2 \int_0^\infty e^{-p^2} dp = 2 \frac{\sqrt{\pi}}{2} = \sqrt{\pi}.\end{aligned}$$

So

$$\int_{-\infty}^\infty S(x, t) dx = \int_{-\infty}^\infty \frac{1}{\sqrt{4\pi kt}} e^{-(x-y)^2/4kt} dx = \frac{1}{\sqrt{\pi}} \int_{-\infty}^\infty e^{-p^2} dp = 1$$

8. We have

$$\max_{\delta \leq x < \infty} S(x, t) = \frac{1}{\sqrt{4\pi kt}} e^{-\delta^2/4kt},$$

so

$$\lim_{t \rightarrow 0} \max_{\delta \leq x < \infty} S(x, t) = \lim_{t \rightarrow 0} \frac{1}{\sqrt{4\pi kt}} e^{-\delta^2/4kt} = \lim_{x \rightarrow +\infty} \frac{\sqrt{x}}{\sqrt{4\pi k}} e^{-x\delta^2/4k} = 0$$

11.(a) Since both  $u(x, t)$  and  $-u(-x, t)$  are solutions, by the uniqueness, we have

$$u(x, t) = -u(-x, t);$$

(b) Since both  $u(x, t)$  and  $-u(-x, t)$  are solutions, by the uniqueness, we have

$$u(x, t) = u(-x, t);$$

(c) similar.

14. Since

$$|e^{-(x-y)^2/4kt} \phi(y)| \leq C e^{-(x-y)^2/4kt + ay^2} = C e^{(a - \frac{1}{4kt})y^2 + \frac{x}{2kt}y - \frac{x^2}{4kt}},$$

so that formula

$$u(x, t) = \frac{1}{\sqrt{4\pi kt}} \int_{-\infty}^\infty e^{-(x-y)^2/4kt} \phi(y) dy$$

makes sense for  $0 < t < 1/(4ak)$ , but not necessarily for larger  $t$ .

**15.** Suppose that  $u, v$  are solution of the diffusion problem with Neumann boundary condition. Let  $w(x, t) = u(x, t) - v(x, t)$ , then

$$w_t = kw_{xx}$$

$$w(x, 0) = w_x(0, t) = w_x(l, t) = 0,$$

and

$$d/dt \int_0^l 1/2[w(x, t)]^2 dx = -k \int_0^l [w_x(x, t)]^2 dx \leq 0.$$

So

$$\int_0^l 1/2[w(x, t)]^2 dx \leq \int_0^l 1/2[w(x, 0)]^2 dx = 0$$

Hence  $w = 0$  and  $u = v$  for all  $t > 0$ .

**16.** Let  $v(x, t) = e^{bt}u(x, t)$ , then

$$v_t - kv_{xx} = 0, \quad v(x, 0) = u(x, 0) = \phi(x)$$

$$\Rightarrow v(x, t) = \frac{1}{\sqrt{4\pi kt}} \int_{-\infty}^{\infty} e^{-(x-y)^2/4kt} \phi(y) dy$$

$$\Rightarrow u(x, t) = \frac{e^{-bt}}{\sqrt{4\pi kt}} \int_{-\infty}^{\infty} e^{-(x-y)^2/4kt} \phi(y) dy$$

**18.** Let  $v(x, t) = u(x + Vt, t)$ , then

$$v_t - kv_{xx} = 0, \quad v(x, 0) = u(x, 0) = \phi(x)$$

$$\Rightarrow v(x, t) = \frac{1}{\sqrt{4\pi kt}} \int_{-\infty}^{\infty} e^{-(x-y)^2/4kt} \phi(y) dy$$

$$\Rightarrow u(x, t) = \frac{1}{\sqrt{4\pi kt}} \int_{-\infty}^{\infty} e^{-(x-Vt-y)^2/4kt} \phi(y) dy$$

**19.** (a)  $S_2(x, y, t) = S(x, t)S(y, t)$ , so

$$S_{2t}(x, y, t) = S_t(x, t)S(y, t) + S(x, t)S_t(y, t)$$

$$S_{2xx}(x, y, t) = S_{xx}(x, t)S(y, t)$$

$$S_{2yy}(x, y, t) = S(x, t)S_{yy}(y, t).$$

Hence

$$\begin{aligned} S_{2t}(x, y, t) &= S_t(x, t)S(y, t) + S(x, t)S_t(y, t) \\ &= kS_{xx}(x, t)S(y, t) + kS(x, t)S_{yy}(y, t) = k(S_{2xx} + S_{2yy}); \end{aligned}$$

(b) We have

$$\begin{aligned} S_2(x, y, t) &= S(x, t)S(y, t) = \\ &= \frac{1}{\sqrt{4\pi kt}} e^{-x^2/4kt} \frac{1}{\sqrt{4\pi kt}} e^{-y^2/4kt} = \frac{1}{4\pi kt} e^{-\frac{x^2+y^2}{4kt}} \end{aligned}$$

Exercise 2.5

1. Let  $u(x, t) = -x^2 - (t - 1)^2$  be the unique solution of the wave equation with boundary conditions:

$$u_{tt} = u_{xx}, \text{ for } -1 < x < 1, 0 < t$$

$$u(x, 0) = -x^2 - 1, \quad u_t(x, 0) = 2, \quad u(-1, t) = u(1, t) = -t^2 + 2t - 2.$$

But  $u(x, t)$  attains its maximum at  $(0, 1)$ .