SHARP ESTIMATES FOR BUBBLING SOLUTIONS OF A FOURTH ORDER MEAN FIELD EQUATION

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ABSTRACT. We consider a sequence of multi-bubble solutions u_k of the following fourth order equation

(*)
$$\Delta^2 u_k = \rho_k \frac{h(x)e^{u_k}}{\int_{\Omega} he^{u_k}} \text{ in } \Omega, \quad u_k = \Delta u_k = 0 \text{ on } \partial\Omega,$$

where h is a $C^{2,\beta}$ positive function, Ω is a bounded and smooth domain in \mathbb{R}^4 , and ρ_k is a constant such that $\rho_k \leq C$. We show that (after extracting a subsequence), $\lim_{k \to +\infty} \rho_k = 32\sigma_3 m$ for some positive integer $m \geq 1$, where σ_3 is the area of the unit sphere in \mathbb{R}^4 . Furthermore, we obtain the following sharp estimates for ρ_k :

$$\rho_k - 32\sigma_3 m = c_0 \sum_{j=1}^m \epsilon_{k,j}^2 (\sum_{l \neq j} \Delta G_4(p_j, p_l) + \Delta R_4(p_j, p_j) + \frac{1}{32\sigma_3} \Delta \log h(p_j)) + o(\sum_{j=1}^m \epsilon_{k,j}^2)$$

where $c_0 > 0$, $\log \frac{64}{\epsilon_{k,j}^4} = \max_{x \in B_{\delta}(p_j)} u_k(x) - \log(\int_{\Omega} h e^{u_k})$ and $u_k \to 32\sigma_3 \sum_{j=1}^m G_4(\cdot, p_j)$ in $C_{loc}^4(\Omega \setminus \{p_1, ..., p_m\})$.

This yields a bound of solutions as ρ_k converges to $32\sigma_3 m$ from below provided that

$$\sum_{i=1}^{m} \left(\sum_{l \neq i} \Delta G_4(p_j, p_l) + \Delta R_4(p_j, p_j) + \frac{1}{32\sigma_3} \Delta \log h(p_j) \right) > 0.$$

The analytic work of this paper is the first step toward computing the Leray-Schauder degree of solutions of equation (*).

1. Introduction

In this paper, we initiate the study of the following fourth order mean field equation

$$\begin{cases} \Delta^2 u = \rho \frac{he^u}{\int_{\Omega} he^u} \text{ in } \Omega, \\ u = \Delta u = 0 \text{ on } \partial\Omega. \end{cases}$$
 (1.1)

This is the first of a series of two papers on computing the Leray-Schauder degree for solutions of (1.1). In this first paper, we compute the sharp estimates of the bubbling rate of multiple bubble solutions.

In dimension two, the analogous problem

$$\begin{cases}
-\Delta u = \rho \frac{he^u}{\int_{\Omega} he^u} \text{ in } \Omega, \\
u = 0 \text{ on } \partial\Omega
\end{cases}$$
(1.2)

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where Ω is a smooth and bounded domain in \mathbb{R}^2 , has been extensively studied by many authors. We summarize the results for (1.2) and identify the difficulty in studying (1.1) now. Let (u_k, ρ_k) be a bubbling sequence to (1.2) with $\rho_k \leq C$, $\max_{x \in \Omega} u_k(x) \to +\infty$. Then it has been proved that

- (P1) (no boundary bubbles) u_k is uniformally bounded near a neighborhood of $\partial\Omega$ (Ma-Wei [18]);
- (P2) (bubbles are simple) $\rho_k \to 8m\pi$ for some $m \ge 1$ and $u_k(x) \to 8\pi \sum_{j=1}^m G_2(\cdot, p_j)$ in $C^2(\Omega \setminus \{p_1, ..., p_m\})$ (Brezis-Merle [2], Nagasaki-Suzuki [21], Li-Shafrir [13], Ma-Wei [18]), where G_2 is the Green function of $-\Delta$ with Dirichlet boundary condition;
- (P3) (sup + inf estimates) at each bubble $p_{k,j}$ where $u_k(p_{k,j}) = \max_{x \in B_\delta(p_j)} u_k(x)$, the following refined estimates hold (Brezis-Li-Shafrir [5], Li [12], Li-Shafrir [13])

$$|u_k(x) - u_k(p_{k,j}) - \log \frac{1}{(1 + \frac{|x - x_{k,j}|^2}{\epsilon_{k,j}^2})^2}| \le C$$
(1.3)

where $u_k(p_{k,j}) - \log(\int_{\Omega} h e^{u_k}) = \log \frac{1}{\epsilon_{k,j}^2}$;

(P4) (exact bubbling rate) It holds then (Chen and Lin [7])

$$\rho_k - 8m\pi = c_0 \sum_{j=1}^m h(p_{k,j})^{-1} \Delta \log h(p_{k,j}) \epsilon_{k,j}^2 \log \frac{1}{\epsilon_{k,j}} + O(\sum_{j=1}^m \epsilon_{k,j}^2);$$
(1.4)

(P5) (**Leray-Schauder degree**) Li [12] initiated the program of computing the Leray-Schauder degree of solutions to (1.2). He showed that the Leray-Schauder degree remains a constant for $\rho \in (8\pi(m-1), 8\pi m)$ and that the degree depends only on the Euler characteristics of Ω . Chen and Lin [8] obtained the exact degree counting formula as follows

$$d(\rho) = \begin{cases} \frac{1}{m!} (-\chi(\Omega) + 1) \cdots (-\chi(\Omega) + m) & \text{for } m > 0, \\ 1 & \text{for } m = 0 \end{cases}$$
 (1.5)

where $\chi(\Omega)$ is the Euler characteristic of Ω .

In this and subsequent paper [17], we carry out the same program for equation (1.1). It will be shown that $d(\rho)$ -the Leray-Schauder degree of (1.1) can be defined as long as $\rho \neq 32m\sigma_3$, where σ_3 is the area of unit sphere in \mathbb{R}^4 . The main purpose of this paper and the subsequent one [17] is to compute $d(\rho)$. In these two papers, we prove, among other things, the following theorem

Theorem A. Let $32m\sigma_3 < \rho < 32(m+1)\sigma_3$ and $d(\rho)$ be the Leray-Schauder degree for equation (1.1). Then

$$d(\rho) = \begin{cases} \frac{1}{m!} (-\chi(\Omega) + 1) \cdots (-\chi(\Omega) + m) & \text{for } m > 0, \\ 1 & \text{for } m = 0 \end{cases}$$
 (1.6)

where $\chi(\Omega)$ is the Euler characteristic of Ω .

Remark. We are informed by Prof. Malchiodi that he obtained a similar degree counting formula for the corresponding prescribing Q-curvature problem on a four dimensional compact manifold, [20]. He used a different approach—the Morse theory approach—to obtain the formula. We remark that on compact manifolds, one doesn't need to prove Property (P1). On the other hand, one of the main difficulties in our proof is the property (P1).

As a consequence of Theorem A, equation (1.1) always possesses a solution for $\rho \neq 32m\sigma_3$ whenever the Euler characteristic $\chi(\Omega) \leq 0$. (Here m can be made ≥ 2 , by results of Lin-Wei [16].) On the other paper, when $\chi(\Omega) > 0$, the situation is much different than the second order case. For example, when Ω is a ball, we can prove the existence of at least one solution when $\rho \in (0, 64\pi\sigma_3)$. See the remark after Corollary (1.2). The complete proof of Theorem A will be given in [17], the second part of this series of papers.

Set $d_m^+ = \lim_{\rho \to 8m\pi^+} d(\rho)$ and $d_m^- = \lim_{\rho \to 32m\sigma_3} d(\rho)$. One of the main steps in the proof of Theorem A is to calculate the gap $d_m^+ - d_m^-$ for any integer $m \ge 1$. Once this is known, $d(\rho)$ can be computed inductively on m. Clearly, the gap of $d_m^+ - d_m^-$ is due to the occurrence of blowup solutions when $\rho \to 32m\sigma_3$. Thus an important question is to analyze the blowup behavior of sequence of solutions u_k to (1.1) and to know the signs $\rho_k - 32m\sigma_3$.

In this paper, we shall obtain estimates analogous to (1.4) for bubbling solutions to (1.1). To this end, we have to first resolve the analogous properties (P1), (P2) and (P3) for problem (1.1). Once we obtain (P1), (P2) follows from results in Wei [23]. So we just need to prove (P1) and (P3). Here the problem arises since the method of Kelvin transform in obtaining (P1) and the method of moving spheres in obtaining (P3) seem not applicable for (1.1). We overcome these difficulties by using various new techniques. (After we obtain (P1)-(P3), the Leray-Schauder degree $d(\rho)$ of (1.1) for $\rho \neq 32m\sigma_3$ can be well-defined.)

The following is the main result of this paper:

Theorem 1.1. Let h be a positive $C^{2,\beta}$ function in Ω and u_k be a sequence of blowup solutions of (1.1) with $\rho = \rho_k$. Then (after extracting a subsequence), $\lim_{k \to +\infty} \rho_k = 32\sigma_3 m$ for some positive integer m. Furthermore,

$$\rho_k - 32\sigma_3 m$$

$$=c_0\sum_{j=1}^{m}(h(p_{k,j}))^{-\frac{1}{2}}\epsilon_{k,j}^2\left[\frac{1}{32\sigma_3}\Delta\log h(p_{k,j})+\Delta R_4(p_{k,j},p_{k,j})+\sum_{i\neq j}\Delta G_4(p_{k,j},p_{k,i})\right]+o(\sum_{j=1}^{m}\epsilon_{k,j}^2)$$

where $c_0 > 0$ is a generic constant, $G_4(\cdot, P)$ is the Green function of Δ^2 with Navier boundary condition $u = \Delta u = 0$ on $\partial\Omega$, R_4 is the regular part of G_4 , $p_{k,j}$ are the local maximum points of u_k on $B_{\delta}(p_j)$, and $\log \frac{64}{\epsilon_{k,j}^4} = u_k(p_{k,j}) - \log(\int_{\Omega} h e^{u_k})$.

Clearly Theorem 1.1 implies the following

Corollary 1.2. Let h(x) be a $C^{2,\beta}$ positive function and satisfy

$$\sum_{j=1}^{m} (h(p_j))^{-\frac{1}{2}} \left[\frac{1}{32\sigma_3} \Delta \log h(p_j) + \Delta R_4(p_j, p_j) + \sum_{l \neq j} \Delta G_4(p_l, p_j) \right] > 0$$
 (1.8)

for all $(p_1, ..., p_m)$ satisfying

$$\nabla \left(\frac{1}{32\sigma_3}\log h(p_j) + R_4(p_j, p_j) + \sum_{l \neq j} G_4(p_l, p_j)\right) = 0, j = 1, ..., m$$
(1.9)

Then for any compact interval $I \subset (32\sigma_3(m-1), 32\sigma_3m]$, there exists a constant C > 0 such that

$$u(x) \le C \text{ for } x \in \Omega \tag{1.10}$$

for any solution u of (1.1) with $\rho \in I$.

As a consequence, if $\Delta h(x) \geq 0$, then (1.10) holds for any solution u of (1.1) with $\rho \in (0, 32\sigma_3]$.

Remark: 1. Corollary 1.2 extends earlier results of Lin and Wei [16] where we proved Corollary 1.2 for m=1, h=1. We note that when $\Omega=B_1$ and h(x)=1, (1.2) has no solution when $\rho \geq 8\pi$. However, for (1.1), a solution exists when $\rho \leq 32\sigma_3$. We conjecture that a solution to (1.1) exists for $any \rho > 0$.

2. Theorem 1.1 can be extended easily to the following n—th order mean field type equation

$$\begin{cases} (-\Delta)^n u = \rho \frac{he^u}{\int_{\Omega} he^u} \text{ in } \Omega, \\ (-\Delta)^j u = 0 \text{ on } \partial\Omega, j = 0, ..., n - 1 \end{cases}$$

$$(1.11)$$

where Ω is a smooth and bounded domain in \mathbb{R}^{2n} . In particular, we have the *same* degree counting formula for solutions to (1.11)

$$d(\rho) = \begin{cases} \frac{1}{m!} (-\chi(\Omega) + 1) \cdots (-\chi(\Omega) + m) & \text{for } m > 0, \\ 1 & \text{for } m = 0 \end{cases}$$

where $\rho \in (m2^{2n}(n-1)!n!\sigma_{2n-1},(m+1)2^{2n}(n-1)!n!\sigma_{2n-1})$. This then implies that (1.11) always has a solution if $\rho \neq m2^{2n}(n-1)!n!\sigma_{2n-1}$ and $\chi(\Omega) \leq 0$.

Semilinear equations involving exponential nonlinearity and fourth order elliptic operator appear naturally in conformal geometry and in particular in prescribing Q-curvature on 4-dimensional Riemannian manifold M (see e.g. Chang-Yang [6])

$$P_g w + 2Q_g = 2\tilde{Q}_{g_w} e^{4w} (1.12)$$

where P_g is the so-called Paneitz operator:

$$P_g = (\Delta_g)^2 + \delta(\frac{2}{3}R_gI - 2\mathrm{Ric}_g)d,$$

 $g_w = e^{2w}g$, Q_g is Q- curvature under the metric g, and \tilde{Q}_{g_w} is the Q-curvature under the new metric g_w .

Integrating (1.12) over M, we obtain

$$k_g := \int_M Q_g = \int_M (\tilde{Q}_{g_w}) e^{4w}$$

where k_g is conformally-invariant. Thus, we can write (1.12) as

$$P_g w + 2Q_g = k_g \frac{\tilde{Q}_{g_w} e^{4w}}{\int_M \tilde{Q}_{g_w} e^{4w}}$$
 (1.13)

In the special case, where the manifold is the Euclidean space, $P_g = \Delta^2$, and (1.13) becomes

$$\Delta^2 w = \rho \frac{h(x)e^{4w}}{\int_{\Omega} h(x)e^{4w}}$$
 (1.14)

There is now an extensive literature about this problem, we refer to Adimurthi-Robert-Struwe [1], Baraket-Dammak-Ouni-Pacard [3], Druet-Robert [9], Hebey-Robert [10], Hebey-Robert-Wen [11], Malchiodi [19] and the references therein.

The organization of this paper is as follows: The statements for properties (P1)-(P3) are collected in Section 2 where important preliminaries are presented. The proof of (P1) is given in the Appendix A and the proof of (P3) is given in Section 3. Finally in Section 4, we prove Theorem 1.1. Though we essentially follow those of [7], we simplify and give a new proof of the key estimates—Estimate C in Section 5.

Throughout this paper, unless otherwise stated, the letter C will always denote various generic constants which are independent of $k \geq 1$.

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2. Preliminaries

We begin with the following lemma which excludes the boundary blowups. The proof of it is by adopting the method used in our previous paper [16] and is given in Appendix A.

Lemma 2.1. Let u be a solution to (1.1) with $\rho \leq C$. Then there exists a $\delta > 0$ such that $u(x) \leq C$ for all x such that $d(x, \partial\Omega) \leq \delta$.

Let G_4 denote the Green's function of Δ^2 under the Navier boundary condition, that is

$$\Delta^2 G_4(x,y) = \delta(x-y), G_4|_{\partial\Omega} = \Delta G_4|_{\partial\Omega} = 0. \tag{2.1}$$

We decompose

$$G_4(x,y) = \frac{1}{4\sigma_3} \log \frac{1}{|x-y|} + R_4(x,y).$$
 (2.2)

It is easy to see that

$$\Delta_x G_4(x, y) < 0, \ \Delta_x R_4(x, y) > 0.$$
 (2.3)

From Lemma 2.1, we derive the following lemma, whose proof follows exactly those in Wei [23] and thus omitted.

Lemma 2.2. Let u_k be a bubbling sequence with $\rho_k \leq C$. Then (after extracting a subsequence), $\rho_k \to 32\sigma_3 m$ and $u_k(x) \to 32\sigma_3 \sum_{j=1}^m G_4(\cdot, p_j)$, where $(p_1, ..., p_m)$ satisfies

$$\nabla \left(\frac{1}{32\sigma_3} \log h(p_i) + R_4(p_i, p_i) + \sum_{j \neq i} G_4(p_i, p_j) \right) = 0, i = 1, ..., m.$$
 (2.4)

We also need to recall the well-known Pohozaev's identity for solutions of fourth-order equation

$$\Delta^2 u = h(x)e^u \text{ in } D.$$

We have

Lemma 2.3. Let u satisfy $\Delta^2 u = h(x)e^u$ in D, where D is a smooth and bounded domain in \mathbb{R}^4 . Then we have

$$\int_{D} (4h + \langle x, \nabla h \rangle) e^{u} = \int_{\partial D} \langle x, \nu \rangle h(x) e^{u}$$

$$+ \int_{\partial D} \left[\frac{1}{2} |\Delta u|^{2} \langle x, \nu \rangle - 2 \frac{\partial u}{\partial \nu} \Delta u - \langle x, \nabla u \rangle \frac{\partial \Delta u}{\partial \nu} - \langle x, \nabla \Delta u \rangle \frac{\partial u}{\partial \nu} + \langle x, \nu \rangle \langle \nabla u, \nabla \Delta u \rangle \right]$$
(2.5)

and for any $\xi \in \mathbb{R}^4$,

$$\int_{D} (\langle \xi, \nabla h \rangle) e^{u} = \int_{\partial D} h(x) e^{u} \langle \xi, \nu \rangle$$
(2.6)

$$\int_{\partial D} \left[\frac{1}{2} |\Delta u|^2 < \xi, \nu > - < \xi, \nabla u > \frac{\partial \Delta u}{\partial \nu} - < \xi, \nabla \Delta u > \frac{\partial \Delta u}{\partial \nu} + < \xi, \nu > < \nabla u, \nabla \Delta u > \right]$$

Proof: In fact, multiplying $\Delta^2 u = h(x)e^u$ by $x \cdot \nabla u$ and integrating by parts, we obtain the lemma.

Let δ_0 be a fixed small constant and $u_k(p_{k,j}) = \max_{x \in B_\delta(p_i)} u_k(x)$ and

$$e^{-c_k} = \frac{1}{\int_{\Omega} h(x)e^{u_k}}$$
 (2.7)

Then $c_k \to +\infty$ as $k \to +\infty$. Let us define

$$l_{k,j} = u_k(p_{k,j}) - c_k, \quad e^{-\frac{l_{k,j}}{4}} = \frac{\epsilon_{k,j}}{\alpha_4^4}, \text{ where } \alpha_4 = 64,$$
 (2.8)

and

$$l_k = \max_{1 \le j \le m} l_{k,j}, \quad \epsilon_k = \min_{1 \le j \le m} \epsilon_{k,j} \tag{2.9}$$

Note that $l_{k,j} \to +\infty$, as otherwise u_k satisfies $|\Delta^2 u_k| \leq C$ in $B_{\delta_0}(p_{k,j})$, $u_k + |\Delta u_k| \leq C$ on $\partial B_{\delta_0}(p_{k,j})$. This implies $\max_{x \in B_{\delta_0}(p_{k,j})} u_k(x) \leq C$, which contradicts to our assumption.

Next, we present a theorem which gives (P3)-sup + inf estimates. The proof of it is interesting and given in a separate section.

Lemma 2.4. We have

$$|u_k(x) - u_k(p_{k,j}) - \log \frac{\alpha_4}{(1 + \frac{|x - p_{k,j}|^2}{\epsilon_k^2})^4}| \le C,$$
 (2.10)

for $x \in B_{\delta_0}(p_{k,j})$.

From Lemma 2.4, we have the following important corollary

Corollary 2.5. Let u_k be a sequence of blowup solutions of (1.1) with $\rho = \rho_k$. Let $l_k, l_{k,j}, \epsilon_k, \epsilon_{k,j}$ be defined as before. It then holds

$$l_k - C \le l_{k,j} \le l_k + C, \ C^{-1} \epsilon_k \le \epsilon_{k,j} \le C \epsilon_k, j = 1, ..., m,$$
 (2.11)

$$c_k - C \le l_{k,j} \le c_k + C, \ C^{-1}e^{-\frac{c_k}{4}} \le \epsilon_{k,j} \le Ce^{-\frac{c_k}{4}}, j = 1, ..., m.$$
 (2.12)

Finally, we consider a problem in \mathbb{R}^4 . It has been proved ([15], [22]) that the solution to the following problem

$$\begin{cases}
\Delta^2 U = e^U, \text{ in } \mathbb{R}^4, \\
\int_{\mathbb{R}^4} e^U < +\infty,
\end{cases}$$
(2.13)

is given by

$$U_{\epsilon,a}(x) := \log \frac{\alpha_4 \epsilon^4}{(\epsilon^2 + |x - a|^2)^4},\tag{2.14}$$

for any $\epsilon > 0$, $a \in \mathbb{R}^4$, provided that

$$U(x) = o(|x|^2) \text{ as } |x| \to +\infty.$$
 (2.15)

Let $U = \log \frac{\alpha_4}{(1+|y|^2)^4}$ and $\tau \in (0,1)$ be a fixed constant. We need the following lemma which proves the nondegeneracy of U:

Lemma 2.6. The solutions to the following linearized problem

$$\Delta^2 \phi = e^U \phi, \quad |\phi(y)| \le C(1+|y|)^{\tau}$$
 (2.16)

is given by $\phi = \sum_{j=0}^4 c_j \psi_j$ where

$$\psi_0 = \frac{1 - |y|^2}{1 + |y|^2}, \psi_j = \frac{y_j}{1 + |y|^2}, j = 1, ..., 4$$
(2.17)

3. Proof of Lemma 2.4

In this section, we prove the **sup+inf estimates**–Lemma 2.4. As we mentioned before, the method of moving spheres seems not applicable here. Instead, we use an approach of combination of potential analysis and Pohozaev identity. This approach has been used in Bartolucci-Chen-Lin-Tarantello [4].

We now state a more general theorem: Let $\tilde{u}_k(x)$ be a solution of

$$\begin{cases} \triangle^2 \tilde{u}_k(x) = h_k(x) e^{\tilde{u}_k} \text{ in } B_2, \text{ and} \\ \int_{B_2} h_k(x) e^{\tilde{u}_k(x)} dx \le C, \end{cases}$$
(3.1)

where $h_k(x)$ converges to a positive function h(x) in $C^1(\bar{B}_2)$, and without loss of generality, we may assume h(0) = 1. Suppose that \tilde{u}_k satisfies the following assumptions

- (i) $|\tilde{u}_k(x) \tilde{u}_k(y)| \le c \text{ for } |x| = |y| = 2,$
- (ii) $|\triangle \tilde{u}_k(x)|$ is bounded in any compact set of $\bar{B}_2 \setminus \{0\}$
- (iii) 0 is the only blow-up point of \tilde{u}_k , i.e., set $S = \{x | x_k \to x \text{ and } \overline{\lim}_{k \to +\infty} \tilde{u}_k(x_k) \to +\infty\}$. Then $S = \{0\}$.

We want to establish the following sharp estimate of the bubbling behavior of \tilde{u}_k near 0. To state our result, we let l_k be the maximum and x_k be a maximum point of \tilde{u}_k , i.e,

$$l_k = \tilde{u}_k(x_k) = \max_{\bar{B}_2} \, \tilde{u}_k \; .$$

and let v(x) be the solution of

$$\Delta^{2}v(x) = e^{v(x)} \text{ in } \mathbb{R}^{4}$$

$$v(0) = 0 = \max_{\mathbb{R}^{2}} v(x) \text{ and } |v(x)| = O(\log|x|) \text{ at } \infty.$$
(3.2)

Theorem 3.1. Suppose \tilde{u}_k is a sequence of solution of (3.1) and satisfies assumptions (i)–(iii) and v is the solution of (3.2). Then there exists a constant c such that

$$|\tilde{u}_k(x) - l_k - v(e^{\frac{l_k}{4}}|x - x_k|)| \le C \text{ in } \bar{B}_1.$$

Applying Theorem 3.1 to $\tilde{u}_k = u_k - c_k$, we obtain Lemma 2.4.

For $r \in (0, 1)$, set

$$\alpha_k(r) = \int_{B_r} h_k(x) e^{\tilde{u}_k(x)} dx ,$$

and

$$\alpha(r) = \lim_{k \to +\infty} \alpha_k(r)$$
 and $\alpha = \lim_{r \to 0} \alpha(r)$.

We first have

Lemma 3.2. Suppose \tilde{u}_k is a solution of (3.1) and satisfies assumption (i)-(iii). Then $\tilde{u}_k \to -\infty$ uniformly in any compact set and $\alpha = 32\sigma_3$.

Proof: Suppose that there exists a point $x_0 \in B_2 \setminus \{0\}$ such that $\tilde{u}_k(x_0)$ is bounded. Then by assumptions (i)-(iii), the sequence \tilde{u}_k is bounded in any compact set of $B_2 \setminus \{0\}$. By taking a diagonal process, a subsequence, still denoted by \tilde{u}_k , approaches a function u(x) in $B_2 \setminus \{0\}$. By Fatou's Lemma,

$$\int_{B_2\setminus\{0\}} e^u < +\infty. \tag{3.3}$$

Since $\int_{B_2} e^{\tilde{u}_k} < +\infty$ and $\tilde{u}_k(x) \to u(x)$ in $B_2 \setminus \{0\}$, we see that $h_k(x)e^{\tilde{u}_k} \to h(x)e^{u(x)} + \alpha'\delta_0$ in the distributional sense, for some constant α' . Here δ_0 is the Dirac measure with singularity at 0. But by (3.3)

$$\alpha' = \lim_{r \to 0} \left(\int_{B_r} h_k(x) e^{\tilde{u}_k} - \int_{B_r} h(x) e^{u(x)} \right) = \alpha.$$

Therefore u(x) satisfies

$$\Delta^2 u(x) = h(x)e^{u(x)} + \alpha \delta_0 \text{ in } B_2.$$

Thus

$$u(x) = \frac{\alpha}{4\sigma_3} \log\left(\frac{1}{|x|}\right) + v(x), \text{ with } \alpha > 0, v(x) \text{ is smooth, and}$$

$$\int_{B_1} e^{u(x)} dx \le C.$$
(3.4)

By the Pohozaev identity (2.5), we have

$$\int_{B_r} [4h_k(x) + (\nabla h_k(x) \cdot x)] e^{\tilde{u}_k}
= \int_{\partial B_r} h_k(x) |x| e^{\tilde{u}_k} d\sigma - \int_{\partial B_r} r \left[\frac{(\Delta \tilde{u}_k)^2}{2} + \frac{\partial \tilde{u}_k}{\partial r} \frac{\partial}{\partial r} \Delta \tilde{u}_k \right] d\sigma
+ \int_{\partial B_r} \frac{\partial}{\partial r} \left(r \frac{\partial \tilde{u}_k}{\partial r} \right) \Delta \tilde{u}_k d\sigma .$$
(3.5)

By letting $k \to +\infty$, we have

$$4\alpha(1+o(1)) = 2\left(\frac{\alpha}{4\sigma_3}\right)^2 \sigma_3(1+o(1)) , \qquad (3.6)$$

where o(1) tends to 0 as $r \to 0$. Since $\alpha > 0$, we have

$$\alpha = 32\sigma_3 \ . \tag{3.7}$$

However, (3.7) implies $\frac{\alpha}{4\sigma_3} = 8$ and then (3.1) yields

$$\int_{B_1} |x|^{-8} dx \le c_1 \int_{B_1} e^{u(x)} dx \le c_1 c ,$$

a contradiction. Thus, $\tilde{u}_k(x) \to -\infty$ uniformly in any compact set of $B_2 \setminus \{0\}$.

Now it is obvious that $\hat{u}_k(x) = \tilde{u}_k(x) - c_k$ converges to $\hat{u}(x)$ in $C^2_{loc}(\bar{B}_2 \setminus \{0\})$ and $c_k \to +\infty$ as $k \to +\infty$ where $c_k = f_{|x|=1}\tilde{u}_k(x)d\sigma$ is the average of \tilde{u}_k over S^3 . Clearly

$$\hat{u}(x) = \frac{\alpha}{4\sigma_3} \log \frac{1}{|x|} + v(x) , \qquad (3.8)$$

with $\triangle^2 v(x) = 0$ in B_2 . We can apply the Pohozaev identity (3.5) to obtain $\alpha = 32\sigma_3$ as the same as (3.7). Thus, Lemma 3.2 is proved.

Proof of Theorem 3.1: By (i) and (ii), it is easy to see that $\tilde{u}_k(x)$ can be written as

$$\tilde{u}_k(x) = \frac{1}{4\sigma_3} \int_{B_2} \log\left(\frac{1}{|x-y|}\right) h_k(y) e^{\tilde{u}_k(y)} dy + f_k(x) , \qquad (3.9)$$

where $f_k(x)$ is a smooth function in $\bar{B}_{\frac{3}{2}}$, and

$$||f_k||_{C^4(\bar{B}_{\frac{3}{2}})} \le C . (3.10)$$

Recall $\tilde{u}_k(x_k) = l_k = \max_{\bar{B}_2} \tilde{u}_k$. Then

$$\tilde{u}_k(x) - l_k = \frac{1}{4\sigma_3} \int_{B_2} \log \left\{ \frac{|x_k - y|}{|x - y|} \right) h_k(y) e^{\tilde{u}_k(y)} dy + f_k(x) - f_k(x_k) . \tag{3.11}$$

Set $v_k(x) = \tilde{u}_k(\varepsilon_k x + x_k) - l_k$ and $\varepsilon_k = e^{-\frac{l_k}{4}}$. Then (3.11) implies

$$v_k(x) = \frac{1}{4\sigma_3} \int_{B_{\varepsilon_k^{-1}}} \log\left(\frac{|y|}{|x-y|}\right) \tilde{h}_k(y) e^{v_k(y)} dy + \tilde{f}_k(x) , \qquad (3.12)$$

where $\tilde{h}_k(y) = h_k(x_k + e^{-\frac{l_k}{4}}y)$ and

$$\|\nabla^j \tilde{f}_k\|_{B_{\varepsilon_k}^{-1}} \to 0 \text{ for } 1 \le j \le 4$$
.

From (3.12), we have $|\triangle v_k(x)|$ is uniformly bounded. Thus, $v_k(x) \to v(x)$ in $C^4_{loc}(\mathbb{R}^4)$ and v(x) satisfies

$$\Delta^2 v(x) = e^{v(x)} \text{ in } \mathbb{R}^4 \text{ and}$$

$$v(x) = \frac{1}{4\sigma_3} \int_{\mathbb{R}^4} \log\left(\frac{|y|}{|x-y|}\right) e^{v(y)} dy + c_0$$
(3.13)

for some constant c_0 . Therefore $\triangle v(x) \to 0$ as $|x| \to +\infty$, a classification result of [2] shows that

$$v(x) = c + \log \frac{1}{(1+|x|^2)^4}$$
.

Thus, for any R > 0,

$$|v_k(x) - v(x)| \to 0$$
 uniformly for $|x| \le R$ (3.14)

as $k \to +\infty$.

To prove Theorem 3.1, it is equivalent to showing

$$|v_k(x) - v(x)| \le C \text{ for } R \le |x| \le r_0 e^{\frac{t_k}{4}},$$
 (3.15)

for some $r_0 > 0$.

To prove (3.15), we claim

$$|\alpha_k - 32\sigma_3| \le c \left(\log \frac{1}{\varepsilon_k}\right)^{-1} , \qquad (3.16)$$

where α_k is the local mass defined by

$$\alpha_k = \int_{B_1} h_k(x) e^{\tilde{u}_k(x)} dx \tag{3.17}$$

The idea to obtain (3.16) is to apply the Pohozaev identity (3.5) on the circle $|x| = \varepsilon_k(\log \frac{1}{\varepsilon_k})$. Hence, we need some fine estimates of v_k . Basically, all estimates required here

can be obtained by using the Green representation formulas (3.12). First, we has a rough estimate about the behavior of v_k .

For any fixed $\delta > 0$, there exists $R = R_{\delta}$ and $k_0 = k(\delta) \in N$ such that if $|x| \geq 2R$ and $k \geq k_0$, then

$$v_k(x) \le -\left(\frac{\alpha_k}{4\sigma_3} - \delta\right) \log|x|$$
.

The proof is standard, and is omitted here. Since $\frac{\alpha_k}{4\sigma_3} \to 8$ as $k \to +\infty$, δ is always chosen such that

$$v_k(x) \le -7\log|x| \quad \text{holds for} \quad |x| \ge 2R.$$
 (3.18)

For $\log \frac{1}{\varepsilon_k} \le |x| \le \varepsilon_k^{-1}$, set

$$\tilde{\alpha}_k(|x|) = \int_{|y| \le r_o|x|} \tilde{h}_k e^{v_k(y)} dy ,$$

where $r_0 \leq \frac{1}{2}$ is a positive constant. By (3.18),

$$|\alpha_k - \tilde{\alpha}_k(|x|)| \le c \int_{|y| \ge r_0|x|} e^{v_k(y)} dy$$

$$\le c \int_{|y| > r_0|x|} |y|^{-7} dy = O\left(\frac{1}{|x|^3}\right),$$
(3.19)

for $|x| \ge \log \frac{1}{\varepsilon_k}$ and k large.

By (3.19), we claim that

$$\left| v_k(x) + \frac{\alpha_k}{4\sigma_3} \log|x| \right| \le C , \qquad (3.20)$$

$$\left| \frac{\partial v_k}{\partial r}(x) + \frac{\alpha_k}{4\sigma_3} \frac{1}{|x|} \right| \le O\left(\left(\log \frac{1}{\varepsilon_k} \right)^{-1} |x|^{-1} \right)$$
 (3.21)

$$\left| \frac{\partial}{\partial r} \left(r \frac{\partial v_k}{\partial r} (x) \right) \right| \le O\left(\frac{1}{|x|^2} \right) \tag{3.22}$$

$$\left| \triangle v_k(x) + \frac{\alpha_k}{2\sigma_3} \frac{1}{|x|^2} \right| \le O\left(\left(\log \frac{1}{\varepsilon_k} \right)^{-1} |x|^{-2} \right) , \text{ and}$$
 (3.23)

$$\left| \frac{\partial}{\partial r} \triangle v_k(x) - \frac{\alpha_k}{\sigma_3} \frac{1}{|x|^3} \right| \le C \left(\log \frac{1}{\varepsilon_k} \right)^{-1} |x|^{-3} , \qquad (3.24)$$

for $|x| = \log \frac{1}{\varepsilon_k}$. In fact, we will prove (3.20) holds for $\log \frac{1}{\varepsilon_k} \le |x| \le \frac{1}{\varepsilon_k}$.

We first show (3.16) by assuming (3.20)-(3.24). Rescaling back to \tilde{u}_k , (3.20)-(3.24) can be written as follows:

$$\tilde{u}_{k}(x) = v_{k} \left(\varepsilon_{k}^{-1}x\right) - 4\log\varepsilon_{k}$$

$$= -\frac{\alpha_{k}}{4\sigma_{3}}\log|x| + \left(\frac{\alpha_{k}}{4\sigma_{3}} - 4\right)\log\varepsilon_{k}$$
(3.25)

$$\frac{\partial \tilde{u}_k(x)}{\partial r} = -\frac{\alpha_k}{4\sigma_3} \frac{1}{|x|} + O\left(\left(\log \frac{1}{\varepsilon_k}\right)^{-1} \frac{1}{|x|}\right) , \qquad (3.26)$$

$$\left| \frac{\partial}{\partial r} \left(r \frac{\partial \tilde{u}_k}{\partial r}(x) \right) \right| \le O\left(\left(\log \frac{1}{\varepsilon_k} \right)^{-1} \frac{1}{|x|} \right) \tag{3.27}$$

$$\Delta \tilde{u}_k(x) = -\frac{\alpha_k}{2\sigma_3} \frac{1}{|x|^2} + O\left(\left(\log \frac{1}{\varepsilon_k}\right)^{-1} \frac{1}{|x|^2}\right), \text{ and}$$
 (3.28)

$$\frac{\partial}{\partial r} \triangle \tilde{u}_k(x) = \frac{\alpha_k}{\sigma_3} \frac{1}{|x|^3} + O\left(\left(\log \frac{1}{\varepsilon_k}\right)^{-1} \frac{1}{|x|^3}\right)$$
(3.29)

for $|x| = 2_k (\log \frac{1}{\varepsilon_k})$.

Substituting (3.25)-(3.29) to (3.5) on $r = \varepsilon_k \log(\frac{1}{\varepsilon_k})$, we have

$$4\alpha_k + O(1)\varepsilon_k \log(\frac{1}{\varepsilon_k}) = \frac{\alpha_k^2}{8\sigma_3} + O((\log \frac{1}{\varepsilon_k})^{-1}).$$

Thus,

$$\alpha_k = 32\sigma_3 + O((\log \frac{1}{\varepsilon_k})^{-1}) ,$$

and then (3.16) is proved.

We come back for the proof of (3.20)-(3.24). By (3.12), for $\log \varepsilon_k^{-1} \leq |x| \leq \varepsilon_k^{-1}$,

$$v_{k}(x) = \frac{1}{4\sigma_{3}} \int_{|y| \leq \varepsilon_{k}^{-1}} \log\left(\frac{|y|}{|x-y|}\right) \tilde{h}_{k}(y) e^{v_{k}(y)} dy + O(1)$$

$$= \frac{1}{4\sigma_{3}} \int_{|y| \leq r_{0}|x|} \log\left(\frac{1}{|x-y|}\right) \tilde{h}_{k}(y) e^{v_{k}(y)} dy + O(1)$$

$$= \frac{\tilde{\alpha}_{k}}{4\sigma_{3}} \log\frac{1}{|x|} + O(1)$$

$$= \frac{\alpha_{k}}{4\sigma_{3}} \log\frac{1}{|x|} + O(1) ,$$

where

$$\begin{split} \int_{|y| \leq e_k^{-1}} |\log(|y|)| \, \tilde{h}_k(y) e^{v_k(y)} dy & \leq c \ , \\ \int_{|y| \geq r_0|x|} \log \left(\frac{1}{|x-y|} \right) \tilde{h}_k(y) e^{v_k(y)} dy & = O(|x|^{-3} \log |x|) \ , \end{split}$$

and (3.19) are employed. This proves (3.20). To prove (3.21), we have

$$\begin{split} \frac{\partial v_k(x)}{\partial r} &= \frac{-1}{4\sigma_3} \int_{|y| \le \varepsilon_k^{-1}} \frac{(x-y) \cdot \frac{x}{|x|}}{|x-y|^2} \tilde{h}_k(y) e^{v_k(y)} dy + O(\varepsilon_k) \\ &= \frac{-1}{4\sigma_3} \int_{|y| \le r_0|x|} \frac{(x-y)x}{|x||x-y|^2} \tilde{h}_k(y) e^{v_k(y)} dy + O(|x|^{-4} + \varepsilon_k) \\ &= \frac{-1}{4\sigma_3} \int_{|y| \le r_0|x|} \frac{1}{|x|} \tilde{h}_k(y) e^{v_k(y)} dy + \frac{O(1)}{|x|^2} \int_{|y| \le r_0|x|} |y| e^{v_k(y)} dy \\ &+ O(|x|^{-4} + \varepsilon_k) \\ &= -\frac{\tilde{\alpha}_k}{4\sigma_3} \frac{1}{|x|} + O(|x|^{-2}) \\ &= -\frac{\alpha_k}{4\sigma_3} |x|^{-1} + O((\log \frac{1}{\varepsilon_k})^{-1} |x|^{-1}). \end{split}$$

For (3.22), we have

$$\begin{split} \frac{\partial}{\partial r} \left(r \frac{\partial v_k}{\partial r}(r) \right) &= \frac{-1}{4\sigma_3} \int_{|y| \le \varepsilon_k^{-1}} \frac{\partial}{\partial r} \left(\frac{(x-y)x}{|x-y|^2} \right) \tilde{h}_k(y) e^{v_k(y)} dy + O(\varepsilon_k) \\ &= \frac{-1}{4\sigma_3} \int_{|y| \le r_0|x|} \frac{\partial}{\partial r} \left(\frac{(x-y)x}{|x-y|^2} \right) \tilde{h}_k(y) e^{v_k(y)} dy + O(\varepsilon_k + |x|^{-4}) \\ &= \frac{O(1)}{|x|^2} \int_{|y| \le r_0|x|} |y| (1+|y|)^{-6} dy + O(\varepsilon_k + |x|^{-4}) \\ &= O\left(\frac{1}{|x|^2}\right) . \end{split}$$

We have proved (3.20)-(3.22). Proofs of (3.23) and (3.24) are similar, and we should omit them here. Hence (3.16) is proved completely.

We note (3.20) holds for $\log \frac{1}{\varepsilon_k} \le |x| \le \varepsilon_k^{-1}$. Therefore, by (3.16) and (3.20), we have

$$|v_k(x) + 8\log|x|| \le C$$

for $\log \frac{1}{\varepsilon_k} \le |x| \le \varepsilon_k^{-1}$. So far, we have proved

$$|v_k(x) - v(x)| \le C \text{ for } |x| \le R \text{ or } |x| \ge \log \frac{1}{\epsilon_k}$$
(3.30)

For the region $R \leq |x| \leq \frac{1}{\varepsilon_k}$, we proceeds as follows: for $|x| \leq \log \frac{1}{\varepsilon_k}$,

$$\begin{split} \Delta v_k(x) &= -\frac{1}{2\sigma_3} \int_{|y| \le r_o|x|} \frac{1}{|x-y|^2} \tilde{h}_k(y) e^{v_k(y)} dy + O(|x|^{-5}) \\ &= -\frac{1}{2\sigma_3} \int_{|y| \le r_o|x|} \frac{1}{|x|^2} \tilde{h}_k(y) e^{v_k(y)} dy - \frac{1}{\sigma_3} \int_{|y| \le r_o|x|} \frac{x \cdot y}{|x|^4} \tilde{h}_k(y) e^{v_k(y)} dy + O(\frac{1}{|x|^4}) \\ &= -\frac{\tilde{\alpha}_k}{2\sigma_3} \frac{1}{|x|^2} + O(\frac{1}{|x|^3}) \\ &= \Delta v(x) + O(\frac{1}{|x|^3}) \end{split}$$

where we have used the fact that $\tilde{\alpha}_k(|x|) - \alpha_k = O(\frac{1}{|x|^3})$ and that $\alpha_k - 32\sigma_3 = O(\frac{1}{\log \frac{1}{\epsilon_k}})$. Thus,

$$|\Delta v_k(x) - \Delta v(x)| \le \frac{C}{|x|^3} \tag{3.31}$$

holds for $R \leq |x| \leq \frac{1}{\epsilon_k}$.

Now we choose $g(x) = C(1 + \frac{1}{|x|})$ where C is large. Then $-\Delta g(x) = \frac{C}{|x|^3} \ge |\Delta v_k(x) - \Delta v(x)|$. It is easy to see that $g(x) \ge |v_k(x) - v(x)|$ for |x| = R and $|x| = \log \frac{1}{\epsilon_k}$. Thus, the maximum principle implies

$$|v_k(x) - v(x)| \le C(1 + \frac{1}{|x|})$$
 (3.32)

for $R \leq |x| \leq \log \frac{1}{\epsilon_k}$.

Combining (3.30) and (3.32), we have proved (3.15). Thus, Theorem 3.1 is completely proved.

4. The Estimate of $\rho_k-32\sigma_4 m$

The main purpose of this section is to prove Theorem 1.1. We follow the main steps used in [7]. Let

$$\tilde{u}_k(x) = u_k(x) - c_k \tag{4.1}$$

which satisfies

$$\Delta^2 \tilde{u}_k = \rho_k h(x) e^{\tilde{u}_k} \text{ in } \Omega, \quad \int_{\Omega} h(x) e^{\tilde{u}_k} = 1.$$
 (4.2)

Recall the definitions $c_k, p_{k,j}, l_k, l_{k,j}, \epsilon_k, \epsilon_{k,j}$ given in (2.7)-(2.9).

For a fixed small $\delta_0 > 0$, we set the local "mass" $\rho_{k,j}$ to be

$$\rho_{k,j} = \rho_k \int_{B_{\delta_0}(p_j)} h(x) e^{\tilde{u}_k(x)} dx \tag{4.3}$$

By Corollary 2.5, we have

$$\rho_{k,j} = \rho_k \int_{B_{\delta_0}(p_{k,j})} h(x) e^{\tilde{u}_k(x)} dx + O(\epsilon_k^4)$$

$$\tag{4.4}$$

which yields

$$\sum_{j=1}^{m} \rho_{k,j} = \rho_k + O(\epsilon_k^4). \tag{4.5}$$

In $B_{\delta_0}(p_{k,j})$, we set

$$G_j^*(x) = \rho_{k,j} R_4(x, p_{k,j}) + \sum_{l \neq j} \rho_{k,l} G_4(x, p_{k,l})$$
(4.6)

and $w_k(x)$ to be the error term defined by

$$w_k(x) = u_k(x) - \sum_{i=1}^m \rho_{k,i} G_4(x, p_{k,i}) = \tilde{u}_k - \sum_{i=1}^m \rho_{k,i} G_4(x, p_{k,i}) + c_k$$
(4.7)

on $\Omega \setminus \bigcup_{j=1}^m B_{\frac{\delta_0}{2}}(p_j)$.

We first have

Estimate A: $|w_k(x)| + |\partial^{\alpha} w_k(x)| = O(\epsilon_k)$ for all $|\alpha| \leq 3, x \in \Omega \setminus \bigcup_{j=1}^m B_{\frac{\delta_0}{2}}(p_j)$.

Proof: This follows from the Green's representation formula:

$$\begin{split} u_k(x) &= \int_{\Omega} G_4(x,y) (\rho_k h(y) e^{\tilde{u}_k}) dy \\ &= \sum_{j=1}^m \rho_k \int_{B_{\frac{\delta_0}{2}}(p_{k,j})} G_4(x,y) h(y) e^{\tilde{u}_k} + O(\epsilon_k^4) \\ &= \sum_{j=1}^m \rho_k \int_{B_{\frac{\delta_0}{2}}(p_{k,j})} [G_4(x,y) - G_4(x,p_{k,j})] h(y) e^{\tilde{u}_k} + \sum_{j=1}^m \rho_k G_4(x,p_{k,j}) \int_{B_{\frac{\delta_0}{2}}(p_{k,j})} h(y) e^{\tilde{u}_k} + O(\epsilon_k^4) \\ &= \sum_{j=1}^m \rho_k \int_{B_{\frac{\delta_0}{2}}(p_{k,j})} O(|y-p_{k,j}|) h(y) e^{\tilde{u}_k} + \sum_{j=1}^m \rho_{k,j} G_4(x,p_{k,j}) + O(\epsilon_k^4) \\ &= \sum_{j=1}^m \rho_{k,j} G_4(x,p_{k,j}) + O(\epsilon_k). \end{split}$$

Similarly we can estimate $|\partial^{\alpha} w_k(x)|$ for $|\alpha| \leq 3$. Hence Estimate A is established.

Estimate B: $|\nabla(\log h(x) + G_j^*(x))| = O(\epsilon_k)$ at $x = p_{k,j}$.

Proof: Applying (2.6) to \tilde{u}_k on $B_{\delta_0}(p_{k,j})$, we obtain

LHS of (2.6) =
$$\rho_k \int_{B_{\delta_0}(p_{k,j})} \langle \xi, \nabla h \rangle e^{\tilde{u}_k}$$

= $\rho_k \int_{B_{\delta_0}(p_{k,j})} [\langle \xi, \nabla h \rangle - \langle \xi, \nabla h(p_{k,j}) \rangle] e^{\tilde{u}_k} + \rho_k \int_{B_{\delta_0}(p_{k,j})} \langle \xi, \nabla h(p_{k,j}) \rangle e^{\tilde{u}_k}$
= $O(\epsilon_k) + \rho_k \int_{B_{\delta_0}(p_{k,j})} h(p_{k,j}) e^{\tilde{u}_k} \langle \xi, \nabla \log h(p_{k,j}) \rangle$
= $O(\epsilon_k) + \rho_{k,j} \langle \xi, \nabla \log h(p_{k,j}) \rangle$.

On the other hand by Estimate A, we have

$$u_k(x) = w_k(x) + \hat{G}_i^*(x) \tag{4.8}$$

where

$$\hat{G}_{j}^{*}(x) = G_{j}^{*}(x) + \frac{\rho_{k,j}}{4\sigma_{3}} \log \frac{1}{|x - p_{k,j}|}.$$

Note that $\Delta^2(\hat{G}_j^*(x)) = 0$ in $B_{\delta_0}(p_{k,j}) \setminus \{p_{k,j}\}$. Applying the Pohozaev's identity to \hat{G}_j^* , we obtain

RHS of (2.6)
$$= \int_{\partial B_{\delta_0}(p_j)} [\langle \xi, \nu \rangle (\frac{1}{2} | \Delta u_k |^2) - \langle \xi, \nabla \Delta u_k \rangle \frac{\partial u_k}{\partial r} - \langle \xi, \nabla u_k \rangle \frac{\partial \Delta u_k}{\partial r} + \langle \xi, \nu \rangle \frac{\partial u_k}{\partial r} \frac{\partial \Delta u_k}{\partial r}] + O(\epsilon_k^4)$$

$$= \int_{\partial B_{\delta_0}(p_j)} [\langle \xi, \nu \rangle (\frac{1}{2} | \Delta \hat{G}_j^* |^2) - \langle \xi, \nabla \Delta \hat{G}_j^* \rangle \frac{\partial \hat{G}_j^*}{\partial r} - \langle \xi, \nabla \hat{G}_j^* \rangle \frac{\partial \Delta \hat{G}_j^*}{\partial r} + \langle \xi, \nu \rangle \frac{\partial \hat{G}_j^*}{\partial r} \frac{\partial \Delta \hat{G}_j^*}{\partial r}]$$

$$+ O(\epsilon_k)$$

$$= \lim_{r \to 0} \int_{\partial B_r(p_j)} [\langle \xi, \nu \rangle (\frac{1}{2} | \Delta \hat{G}_j^* |^2) - \langle \xi, \nabla \Delta \hat{G}_j^* \rangle \frac{\partial \hat{G}_j^*}{\partial r} - \langle \xi, \nabla \hat{G}_j^* \rangle \frac{\partial \Delta \hat{G}_j^*}{\partial r} + \langle \xi, \nu \rangle \frac{\partial \hat{G}_j^*}{\partial r} \frac{\partial \Delta \hat{G}_j^*}{\partial r}]$$

$$+ O(\epsilon_k)$$

$$= -\rho_{k,j} \langle \xi, \nabla G_j^* (p_{k,j}) \rangle + O(\epsilon_k)$$

This proves Estimate B.

Next, we give a sharper description of the bubbling behavior of u_k in the ball $B_{\delta_0}(p_{k,j})$. We set

$$h_{k,j} = h(p_{k,j})$$

and

$$v_{k,j}(x) = \log\left(\frac{\alpha_4 \epsilon_{k,j}^4}{(\epsilon_{k,j}^2 + \sqrt{\rho_k h_{k,j}} |x - q_{k,j}|^2)^4}\right)$$
(4.9)

where $\epsilon_{k,j}$ is given by (2.8) and $q_{k,j} \in \mathbb{R}^4$ is chosen such that

$$\nabla v_{k,j}(p_{k,j}) = \nabla \log h(p_{k,j}). \tag{4.10}$$

By direct computations, we also have

$$|p_{k,j} - q_{k,j}| = O(\epsilon_k^2) \tag{4.11}$$

For $x \in B_{\delta_0}(p_{k,j})$, we also set

$$\eta_{k,j}(x) = u_k(x) - c_k - v_{k,j}(x) - (G_j^*(x) - G_j^*(p_{k,j})). \tag{4.12}$$

Then by Estimate B and (4.11), we have

$$|\nabla \eta_{k,j}(p_{k,j})| = O(\epsilon_k) \tag{4.13}$$

$$\eta_{k,j}(p_{k,j}) = l_{k,j} - v_{k,j}(p_{k,j}) = O(\epsilon_k)$$
(4.14)

$$|\eta_{k,j}(x)| \le C \quad \text{for } x \in B_{\delta_0}(p_{k,j}) \tag{4.15}$$

Let us check the matching of w_k and $\eta_{k,j}$ on $B_{\frac{\delta}{2}}(p_{k,j})$:

$$\eta_{k,j}(x) = u_k(x) - c_k - v_{k,j}(x) - (G_j^*(x) - G_j^*(p_{k,j}))$$

$$= w_k + \sum_i \rho_{k,i} G_4(x, p_{k,i}) - c_k - \log \frac{\alpha_4 \epsilon_j^4}{(\epsilon_j^2 + \sqrt{\rho_k h_{k,j}} | x - q_{k,j}|^2)^4} - (G_j^*(x) - G_j^*(p_{k,j}))$$

$$= O(\epsilon_k) + \rho_{k,j} G_4(x, p_{k,j}) + \rho_{k,j} R_4(x, p_{k,j}) - c_k + G_j^*(p_{k,j}) - \log \frac{\alpha_4 \epsilon_{k,j}^4}{(\sqrt{\rho_k h_{k,j}} | x - q_{k,j}|^2)^4}$$

$$= O(\epsilon_k) + (\frac{\rho_{k,j}}{4\sigma_3} - 8) \log \frac{1}{|x - p_{k,j}|} - c_k + G_j^*(p_{k,j}) - \log \frac{\alpha_4 \epsilon_{k,j}^4}{(\sqrt{\rho_k h_{k,j}})^4}$$

$$= O(\epsilon_k) + A_{k,j} + (\frac{\rho_{k,j}}{4\sigma_2} - 8) \log \frac{1}{|x - p_{k,j}|}$$

where $A_{k,j}$ is a constant given by

$$A_{k,j} = -c_k - \log \epsilon_j^4 + G_j^*(P_{k,j}) - \log \frac{\alpha_4}{\rho_k^2 h_{k,j}^2}.$$
 (4.16)

Then by Corollary 2.5, we have $|A_{k,j}| = O(1)$. This implies that for $x \in \partial B_{\delta_0}(p_{k,j})$, we have

$$\eta_{k,j}(x) = A_{k,j} + (\frac{\rho_{k,j}}{4\sigma_3} - 8)\log\frac{1}{|x - p_{k,j}|} + O(\epsilon_k). \tag{4.17}$$

Moreover, (4.17) holds for partial derivatives of $\eta_{k,j}$ up to the order 3.

Let

$$\tilde{\eta}_{k,j}(z) = \eta_{k,j}(p_{k,j} + \epsilon_{k,j}(\rho_k h_{k,j})^{-1/4} z), \quad R = \frac{\delta_0}{8\epsilon_{k,j}}.$$
 (4.18)

The following estimate is the key estimate, whose proof will be given in a separate section

Estimate C: For any $\tau \in (0,1)$, there exists a constant C_{τ} such that

$$|\tilde{\eta}_{k,j}(z)| \le C_{\tau} (1+|z|)^{\tau} (\epsilon^{2\tau} + \epsilon^{\tau} \sup_{\frac{R}{2} \le |z| \le R} |\tilde{\eta}_{k,j}(z)|)$$
 (4.19)

Now we have

Estimate D: For any $\frac{1}{2} < \tau < 1$, we have

$$|A_{k,j}| \le C(\epsilon + \epsilon^{\tau} \sup_{\frac{R}{2} \le |z| \le R} |\tilde{\eta}_{k,j}|) \tag{4.20}$$

Proof: By Green's formula,

$$\begin{array}{lcl} u_k(p_{k,j}) & = & \log(\alpha_4 \epsilon_{k,j}^{-4}) + c_k \\ \\ & = & \sum_{i=1}^m \int_{B_{\delta_0}(p_{k,l})} \rho_k h e^{\tilde{u}_k} G_4(p_{k,j},y) dy + O(\epsilon_k). \end{array}$$

For $l \neq j$,

$$\int_{B_{\delta_0}(p_{k,l})} \rho_k h e^{\tilde{u}_k} G_4(p_{k,j}, y) dy = \rho_{k,l} G_4(p_{k,j}, p_{k,l}) + O(\epsilon_k).$$
(4.21)

For l = j, we have

$$\begin{split} & \int_{B_{\delta_0}(p_{k,j})} \rho_k h e^{\tilde{u}_k} G_4(p_{k,j},y) dy \\ = & \int_{B_{\delta_0}(p_{k,j})} \rho_k h e^{\tilde{u}_k} (\frac{1}{4\sigma_3} \log \frac{1}{|p_{k,j}-y|} + R_4(p_{k,j},y)) dy \\ = & -\frac{1}{4\sigma_3} \rho_k \int_{B_{\delta_0}(p_{k,j})} \rho_k h e^{\tilde{u}_k} \log \frac{1}{|p_{k,j}-y|} dy + \rho_{k,j} R_4(p_{k,j},p_{k,j}) + O(\epsilon_k). \end{split}$$

Now we write

$$he^{\tilde{u}_k} = he^{v_{k,j} + \eta_{k,j} + (G_j^*(x) - G_j^*(p_{k,j}))}$$

= $h(p_{k,j})e^{v_{k,j}}H(x, \eta_{k,j}(x)) + h(p_{k,j})e^{v_{k,j}}$

where

$$H(x,t) = e^{t + \log h(x) + G_j^*(x) - \log h(p_{k,j}) - G_j^*(p_{k,j})} - 1$$
(4.22)

Let z and $z_{k,j}$ satisfy

$$x = p_{k,j} + \epsilon_{k,j} (\rho_k h_{k,j})^{-1/4} z, \quad q_{k,j} = p_{k,j} + \epsilon_{k,j} (\rho_k h_{k,j})^{-1/4} z_{k,j}$$

Then we have

$$\rho_{k} \int_{B_{\delta}(P_{k,j})} he^{\tilde{u}_{k}} \log |p_{k,j} - y|
= \rho_{k} \int_{B_{\delta_{0}}(p_{k,j})} he^{\tilde{u}_{k}} \log \epsilon_{k,j} (\rho_{k} h_{k,j})^{-1/4} + \rho_{k} \int_{B_{R}(0)} he^{\tilde{u}_{k}} \log |z|
= \rho_{k,j} \log(\epsilon_{k,j} (\rho_{k} h_{k,j})^{-1/4}) + \rho_{k} \int_{B_{R}(0)} \frac{\alpha_{4}}{(1 + |z - z_{k,j}|^{2})^{4}} (1 + H(x, \eta_{k,j})) \log |z|
= \rho_{k,j} \log(\epsilon_{k,j} (\rho_{k} h_{k,j})^{-1/4}) + \rho_{k} \int_{B_{R}(0)} \frac{\alpha_{4}}{(1 + |z - z_{k,j}|^{2})^{4}} H(x, \eta_{k,j}) \log |z| + O(\epsilon_{k}) (4.23)$$

where we have used

$$|z_{k,j}| = O(\epsilon_k), \int_0^\infty (1+r^2)^{-4} r^3 \log r dr = 0.$$
 (4.24)

The last term in (4.23) can be estimated by

$$\int_{B_{R}(0)} \frac{\alpha_{4}}{(1+|z-z_{k,j}|^{2})^{4}} H(x,0) \log|z| + O(\epsilon_{k}^{2\tau} + \epsilon_{k}^{\tau} \sup_{\frac{R}{2} \leq |z| \leq R} |\tilde{\eta}_{k,j}(z)|)
= O(\epsilon_{k} + \epsilon_{k}^{\tau} \sup_{\frac{R}{2} \leq |z| \leq R} |\tilde{\eta}_{k,j}(z)|).$$

Combining all together, we obtain Estimate D (by choosing a larger $\tau > \frac{1}{2}$). Estimate D implies

Estimate E: On $B_{\frac{\delta_0}{2}}(p_{k,j})$

$$\eta_{k,j}(x) = \left(\frac{\rho_{k,j}}{4\sigma_3} - 8\right) \log \frac{1}{|x - p_{k,j}|} + O(\epsilon + \epsilon^{\tau} \sup_{\frac{R}{\alpha} < |z| < R} |\tilde{\eta}_{k,j}(z)|). \tag{4.25}$$

From (4.25), we also have

$$\sup_{\frac{R}{3} < |z| < R} |\tilde{\eta}_{k,j}| \le C|\rho_{k,j} - 32\sigma_3| + C\epsilon_{k,j}. \tag{4.26}$$

Hence Estimate C can be refined as

$$|\tilde{\eta}_{k,j}(z)| \le C(1+|z|)^{\tau} (\epsilon_k^{2\tau} + \epsilon_k^{\tau} |\rho_{k,j} - 32\sigma_3|).$$
 (4.27)

On the other hand, $|\rho_{k,j}-32\sigma_3|$ can also be estimated by a quantity related to $\eta_{k,j}$.

Estimate F

$$\rho_{k,j} - 32\sigma_3 = \int_{\partial B_{\delta_0}(p_{k,j})} \frac{\partial \Delta \eta_{k,j}}{\partial \nu} dx + o(\epsilon_k^2). \tag{4.28}$$

Proof:

$$\rho_{k,j} = \rho_k \int_{B_{\delta}(p_{k,j})} h e^{\tilde{u}_k} = \int_{\partial B_{\delta_0}(p_{k,j})} \frac{\partial \Delta u_k}{\partial \nu}$$

$$= \int_{\partial B_{\delta_0}(p_{k,j})} \frac{\partial \Delta v_{k,j}}{\partial \nu} + \int_{\partial B_{\delta_0}(p_{k,j})} \frac{\partial \Delta \eta_{k,j}}{\partial \nu}$$

$$= 32\sigma_3 + \int_{\partial B_{\delta_0}(p_{k,j})} \frac{\partial \Delta \eta_{k,j}}{\partial \nu} + O(\epsilon_k^3)$$

since

$$\frac{\partial \Delta v_{k,j}}{\partial \nu} = 32 \frac{1}{r^3} + O(\frac{1}{r^7}).$$

It remains to compute $\int_{\partial B_{\delta_0}(p_{k,j})} \frac{\partial \Delta \eta_{k,j}}{\partial \nu}$.

Note that $\eta_{k,j}$ satisfies

$$\Delta^2 \eta_{k,j} = \rho_k h_{k,j} e^{v_{k,j}} \eta_{k,j} + \rho_k h_{k,j} e^{v_{k,j}} (H(x, \eta_{k,j}) - \eta_{k,j}). \tag{4.29}$$

Let

$$\Psi(x) = \frac{\epsilon_{k,j}^2 - \sqrt{\rho_k h_{k,j}} |x - q_{k,j}|^2}{\epsilon_{k,j}^2 + \sqrt{\rho_k h_{k,j}} |x - q_{k,j}|^2}.$$

Then it is easy to see that $\Psi(x)$ satisfies

$$\Delta^2 \Psi(x) = \rho_k h_{k,j} e^{v_{k,j}(x)} \Psi(x). \tag{4.30}$$

Using (4.29) and (4.30), we obtain

$$\int_{B_{\delta_0}(p_{k,j})} [\Delta^2 \eta_{k,j} \Psi(x) - \Delta^2 \Psi(x) \eta_{k,j}(x)] = \int_{B_{\delta_0}(p_{k,j})} \rho_k h_{k,j} e^{v_{k,j}} [H(x, \eta_{k,j}) - \eta_{k,j}] \Psi(x). \tag{4.31}$$

The left-hand-side of (4.31) equals

$$= \int_{\partial B_{\delta_0}(p_{k,j})} \left[\frac{\partial \Delta \eta_{k,j}}{\partial \nu} \Psi(x) - \Delta \eta_{k,j} \frac{\partial \Psi}{\partial \nu} + \Delta \Psi \frac{\partial \eta_{k,j}}{\partial \nu} - \eta_{k,j} \frac{\partial \Delta \Psi}{\partial \nu} \right]$$

$$= \int_{\partial B_{\delta_0}(p_{k,j})} \frac{\partial \Delta \eta_{k,j}}{\partial \nu} (\Psi(x) + 1) - \int_{\partial B_{\delta_0}(p_{k,j})} \frac{\partial \Delta \eta_{k,j}}{\partial \nu} + O(\epsilon_{k,j}^2 | \rho_{k,j} - 32\sigma_3| + \epsilon_{k,j}^3)$$

$$= \int_{\partial B_{\delta_0}(p_{k,j})} \frac{\partial \Delta \eta_{k,j}}{\partial \nu} \frac{2\epsilon_{k,j}^2}{\epsilon_{k,j}^2 + \sqrt{\rho_k h_{k,j}}} \delta_0^2 + o(\epsilon_{k,j}^2) - \int_{\partial B_{\delta_0}(p_{k,j})} \frac{\partial \Delta \eta_{k,j}}{\partial \nu} + O(\epsilon_{k,j}^2 | \rho_{k,j} - 32\sigma_3| + \epsilon_{k,j}^3)$$

$$= -\int_{\partial B_{\delta_0}(p_{k,j})} \frac{\partial \Delta \eta_{k,j}}{\partial \nu} + o(\epsilon_{k,j}^2).$$

The right-hand-side of (4.31) equals

$$=\frac{1}{\sqrt{\rho_k h_{k,j}}}\int_{\mathbb{R}^4} (\frac{1}{(1+|z-z_{k,j}|^2)^4}) \left[\epsilon_{k,j}^2 \sum_{l,m} \frac{\partial^2}{\partial x_l \partial x_m} [G_j^* + \log h] z_l z_m + O(|\eta_{k,j}|^2) \right] \frac{1-|z-z_{k,j}|^2}{1+|z-z_{k,j}|^2} dz$$

$$= \frac{1}{\sqrt{\rho_k h_{k,j}}} \int_{\mathbb{R}^4} \left(\frac{1}{(1+|z|^2)^4}\right) \left[\epsilon_{k,j}^2 \sum_{l,m} \frac{\partial^2}{\partial x_l \partial x_m} [G_j^* + \log h](z_l + z_{k,l})(z_m + z_{k,m}) + O(\epsilon_k^{4\tau})\right] \frac{1 - |z|^2}{1 + |z|^2} dz$$

$$= \frac{1}{\sqrt{\rho_k h(p_{k,j})}} \epsilon_{k,j}^2 \Delta(G_j^* + \log h)(p_{k,j}) \frac{1}{4} \int_{\mathbb{R}^4} \frac{|z|^2}{(1+|z|^2)^4} \frac{1-|z|^2}{1+|z|^2} dz + o(\epsilon_{k,j}^2)$$
(4.32)

where

$$\frac{1}{4} \int_{\mathbb{R}^4} \frac{|z|^2}{(1+|z|^2)^4} \frac{1-|z|^2}{1+|z|^2} dz = \frac{1}{4} \int_{\mathbb{R}^4} \frac{|z|^2-1}{(1+|z|^2)^4} \frac{1-|z|^2}{1+|z|^2} dz < 0. \tag{4.33}$$

Combining Estimate F and (4.32), we obtain

$$\rho_{k,j} - 32\sigma_4 = c_0 \frac{1}{\sqrt{h(p_{k,j})}} \epsilon_{k,j}^2 \Delta(G_j^*(p_{k,j}) + \log h(p_{k,j})) + o(\epsilon_{k,j}^2)$$
(4.34)

where $c_0 > 0$.

Finally, summing up the estimates in (4.34), we obtain Theorem 1.1.

5. Proof of Estimate C

In this section, we prove Estimate C. Our proof is different from [7] and is simpler.

Let $\tau \in (0,1)$ be a fixed positive number. We begin with the following simple but important lemma

Lemma 5.1. Let u satisfy

$$\Delta^2 u = f(y)$$
 in $B_R(0)$, $u = \Delta u = 0$ on $\partial B_R(0)$

Then for R large we have

$$\| \langle y \rangle^{-\tau} (u - u(0)) \|_{L^{\infty}(B_{\frac{R}{2}}(0))} \le C \| \langle y \rangle^{4-\tau} f(y) \|_{L^{\infty}(B_{R}(0))}.$$
 (5.1)

Proof: Assume that $|| < y >^{4-\tau} f(y)||_{L^{\infty}(B_R(0))} = 1$.

By the Green's representation formula,

$$u(y) - u(0) = \int_{B_R(0)} \left[\frac{1}{4\sigma_3} \log \frac{|z|}{|y - z|} + \hat{R}_4(\frac{y}{R}, \frac{z}{R}) - \hat{R}_4(0, \frac{z}{R}) \right] f(z) dz$$
 (5.2)

where \hat{R}_4 is the regular part of the Green's function of Δ^2 in $B_1(0)$ with Navier boundary condition.

Note that

$$\int_{B_R(0)} |\hat{R}_4(\frac{y}{R}, \frac{z}{R}) - \hat{R}_4(0, \frac{z}{R})||f(z)|dz \le C \frac{1}{R^\tau} \int_{B_R(0)} \langle y \rangle^\tau |f(z)|dz$$

$$\leq C < y >^{\tau}$$
.

So we only need to consider the first term on the right hand side of (5.2). To this end, we decompose

$$\int_{B_R(0)} \log \frac{|z|}{|y-z|} f(z) dz = \left(\int_{|z| \le \frac{1}{2}|y|} + \int_{\frac{1}{2}|y| \le |z| \le 2|y|} + \int_{|z| \ge 2|y|} \right) \log \frac{|z|}{|y-z|} f(z) dz.$$

The last integral can be controlled by

$$|\int_{|z|>2|y|} \log \frac{|z|}{|y-z|} f(z) dz| \le C |\int_{|z|>2|y|} \frac{|y|}{|z|} |f(z)| dz \le C < y >^{\tau}.$$

For the first integral, we have

$$|y - z| \ge |y| - |z| \ge \frac{1}{2}|y| \ge |z|$$

$$|\int_{|z| \le \frac{1}{2}|y|} \log \frac{|z|}{|y - z|} f(z) dz| \le \int_{|z| \le \frac{1}{2}|y|} \log \frac{|y - z|}{|z|} < z >^{\tau - 4} dz$$

$$\le C < y >^{\tau}.$$

It remains to compute the last integral

$$\int_{\frac{1}{2}|y| \le |z| \le 2|y|} \log \frac{|z|}{|y-z|} f(z) dz = |y|^4 \int_{\frac{1}{2} \le |\tilde{z}| \le 2} \log \frac{|\tilde{z}|}{|e_1 - \tilde{z}|} f(|y|\tilde{z}) d\tilde{z}
\le |y|^4 \int_{\frac{1}{2} \le |\tilde{z}| \le 2} |\log \frac{|\tilde{z}|}{|e_1 - \tilde{z}|} | < y >^{\tau - 4} < \tilde{z} >^{\tau - 4} d\tilde{z}
< C < y >^{\tau}.$$

Let us go back to the equation for $\tilde{\eta}$:

$$\begin{split} \Delta^2 \tilde{\eta} &= e^U(e^{\tilde{\eta}} - 1) + O(\epsilon_{k,j}^2 < y >^{-6}) &\text{in } B_{\frac{\delta_0}{\epsilon_{k,j}}}(0) \\ \tilde{\eta} &= O(1), \Delta \tilde{\eta} = O(\epsilon_{k,j}^2) &\text{on } \partial B_{\frac{\delta_0}{\epsilon_{k,j}}}(0). \end{split}$$

Let $\hat{\eta}_{k,j} = \tilde{\eta}_{k,j} \chi$. Then we have $\hat{\eta}_{k,j}$ satisfies

$$\Delta^{2}\hat{\eta}_{k,j} = e^{U} \frac{(e^{\tilde{\eta}_{k,j}} - 1)}{\tilde{\eta}_{k,j}} \hat{\eta}_{k,j} + O(\epsilon_{k,j}^{2} < y >^{-6}) \text{ in } B_{\frac{\delta_{0}}{\epsilon_{k,j}}}(0), \quad \hat{\eta}_{k,j} = \Delta \hat{\eta}_{k,j} = 0 \text{ on } \partial B_{\frac{\delta_{0}}{\epsilon_{k,j}}}(0).$$
(5.3)

We claim that

$$\| \langle y \rangle^{-\tau} \hat{\eta}_{k,j} \|_{L^{\infty}(B_{\frac{\delta}{2\epsilon_{k,j}}}(0))} \le C\epsilon_{k,j}^{\tau}.$$
 (5.4)

Estimate C follows from (5.4).

Suppose not. Then $\epsilon_{k,j}^{-\tau} \| < y >^{-\tau} \hat{\eta}_{k,j} \|_{L^{\infty}(B_{\frac{\delta_0}{2\epsilon_{k,j}}}(0))} \to +\infty$. Let $\bar{\eta}_{k,j} = \frac{\hat{\eta}_{k,j}}{\| < y >^{-\tau} \hat{\eta}_{k,j} \|_{L^{\infty}(B_{\frac{\delta}{2\epsilon_{k,j}}}(0))}}$.

Then $\bar{\eta}_{k,j}$ satisfies

$$\Delta^2 \bar{\eta}_{k,j} = e^U \frac{(e^{\tilde{\eta}_{k,j}} - 1)}{\tilde{\eta}_{k,j}} \bar{\eta}_{k,j} + o(\langle y \rangle^{\tau-4}) \text{ in } B_{\frac{\delta}{\epsilon_{k,j}}}(0), \quad \bar{\eta}_{k,j} = \Delta \bar{\eta}_{k,j} = O \text{ on } \partial B_{\frac{\delta}{\epsilon_{k,j}}}(0)$$

where $\bar{\eta}_{k,j}(0) = 0$.

We claim that $\bar{\eta}_{k,j} \to 0$ in $C^1_{loc}(\mathbb{R}^4)$. In fact, by standard elliptic regularity theory, $\bar{\eta}_{k,j} \to \eta_0$, where η_0 satisfies

$$\Delta^2 \eta_0 = e^U \eta_0, \eta_0(0) = \nabla \eta_0(0) = 0, |\eta_0(y)| \le C < y >^{\tau}.$$
(5.5)

By Lemma 2.6, $\eta_0 = \sum_{j=0}^4 c_j \psi_j$ for some constants c_j , j = 0, 1, ..., 4. Using the assumption $\eta_0(0) = \nabla \eta_0(0) = 0$, we deduce that $c_j = 0$ and hence $\eta_0 \equiv 0$.

So $\bar{\eta}_{k,j} \to 0$ in $C^1_{loc}(\mathbb{R}^4)$. Now we consider

$$\| < y >^{4-\tau} e^{U} \frac{(e^{\bar{\eta}_{k,j}} - 1)}{\tilde{\eta}_{k,j}} \bar{\eta}_{k,j} \|_{L^{\infty}(B_{\frac{\delta_0}{\epsilon_{k,j}}}(0))} \le C \| < y >^{4-\tau} < y >^{-8} \bar{\eta}_{k,j} \|_{L^{\infty}(B_{\frac{\delta}{\epsilon_{k,j}}}(0))} = o(1).$$

By Lemma 5.1, we conclude

$$\| < y >^{-\tau} \bar{\eta}_{k,j} \|_{L^{\infty}(B_{\frac{\delta}{2\epsilon_{k,j}}}(0))} \le C \| < y >^{-\tau-4} e^{U}(e^{\tilde{\eta}_{k,j}} - 1) \tilde{\eta}_{k,j}^{-1} \bar{\eta}_{k,j} \|_{L^{\infty}(B_{\frac{\delta}{2\epsilon_{k,j}}}(0))} + o(1) = o(1)$$

which contradicts to the assumption that $\| < y >^{-\tau} \bar{\eta}_{k,j} \|_{L^{\infty}(B_{\frac{\delta}{2\epsilon_{k,j}}}(0))} = 1$.

Appendix A: Proof of Lemma 2.1

In this appendix, we prove Lemma 2.1. We follow the ideas in Section 2 of our previous paper [16].

Let u_k be a sequence of solutions of

$$\begin{cases} \Delta^2 u_k = \mu_k h(x) e^{u_k}, & \text{in } Om \\ u_k = 0 & \text{on } \partial \Omega, \end{cases}$$

and u_k blows up at $\{q_1, \ldots, q_n\}$, where $\mu_k = \frac{\rho_k}{\int_{\Omega} h(x)e^{u_k}} \to 0$.

In our previous paper [16], we have shown that $u_k \to P(x)$ in $C^4(\bar{\Omega} \setminus \{p_1, \ldots, p_m\})$. It is not difficult to show that

$$\lim_{x \to p_j} P(x) = \lim_{x \to p_j} (-\Delta P(x)) = +\infty, j = 1, ..., m$$

Claim. $p_j \in \Omega \text{ for } j = 1, 2, \dots, m.$

We prove it by contradiction. Assume $p_1 = 0 \in \partial\Omega$ and $e_1 = (1, 0, ..., 0)$ is the outernormal of $\partial\Omega$ at zero. (See Figure 2.1 in [16].)

Let $N = B(0, r) \cap \Omega$ and φ be a solution of

$$\begin{cases} \Delta^{2}\varphi = 0 & \text{in } \Omega, \\ \Delta\varphi < 0 & \text{and } \varphi > 0 & \text{in } \Omega, \\ \varphi = 0 & \text{on } B(0, r) \cap \partial\Omega, \\ \Delta\varphi = 0 & \text{on } B(0, r) \cap \partial\Omega \end{cases}$$

We now choose C is a large number so that

$$he^{C\varphi(x)}$$
 is decreasing in x_1 for $x \in N$.

N can be chosen so small such that

$$P(x) - C\varphi(x) \ge 0$$
 and $\Delta P(x) - C\Delta\varphi(x) < 0$

for $x \in \partial N$. Therefore by the maximum principle, for large k,

$$u_k - C\varphi(x) > 0$$
 and $\Delta u_k(x) - C\Delta\varphi(x) < 0$

for $x \in \partial N \cap \Omega$. Set $w_k = u_k - C\varphi(x)$. Then

$$\Delta^2 w_k = \mu_k (he^{C\varphi}) e^{w_k} \text{ in } N$$

$$w_k = \Delta w_k = 0 \text{ for } x \in \partial \Omega \cap N.$$

Then we can use the local version of the moving planes as before to show $w_k(x)$ is decreasing in x_1 , which yields a contradiction to the assumption 0 is a blowup point.

In this section, we prove the nondegeneracy Lemma 2.6

Let $U = \log \frac{\alpha_4}{(1+|y|^2)^4}$ be a solution of $\Delta^2 u = e^u$. In this section, we study the following eigenvalue problem

$$\Delta^2 \phi = e^U \phi, \quad |\phi| \le C < y >^{\tau} \tag{5.6}$$

for some $\tau \in (0,1)$.

Since ψ_j , j=0,...,4 satisfies (5.6), we may assume that ϕ satisfies $\int_{\mathbb{R}^4} e^U \psi_j \phi = 0$, j=0,1,...,4. Our goal is then to show that $\phi \equiv 0$.

Let $\phi_0 = \frac{1}{4\sigma_3} \int_{\mathbb{R}^4} \log \frac{1}{|y-z|} e^U \phi dz$. Then we have $\Delta^2(\phi - \phi_0) = 0$. Since $|\phi_0(y)| \leq C \log(2+|y|)$, we have $\phi - \phi_0 = C$. Thus,

$$\phi(y) = \frac{1}{4\sigma_3} \int_{\mathbb{R}^4} \log \frac{1}{|y - z|} e^U \phi dz + C.$$
 (5.7)

We first prove

Lemma 5.2. Let ϕ satisfy (5.6). Then $|\phi| \leq C$ and there holds

$$\int_{\mathbb{R}^4} e^U \phi = 0, \int_{\mathbb{R}^4} e^U \psi_j \phi = 0, j = 0, 1, ..., 4.$$
 (5.8)

Proof: From (5.7), we see that

$$\phi = \ln|y| \left(\frac{1}{4\sigma_3} \int_{\mathbb{R}^4} e^U \phi\right) + O(1), \text{ for } |y| \ge 1$$
 (5.9)

and (5.9) holds for $D^{\alpha}\phi$, $|\alpha| \leq 3$.

We decompose

$$\phi = \phi_0(r) + \phi' \tag{5.10}$$

where $\phi_0(r) = \frac{1}{|S^3|} \int_{S^3} \phi d\sigma$. Then it is easy to see that both ϕ_0 and ϕ' satisfy (5.6). Since U is radially symmetric, by (5.9), it is easy to see that ϕ' is bounded and satisfies (5.8). Therefore, to prove the Lemma, it is enough to assume that $\phi = \phi(r)$ is radially symmetric. Now multiplying (5.6) by $\psi_0(r) = \frac{1-r^2}{1+r^2}$ and integrating over $B_r(0)$, we obtain

$$0 = \int_{\partial B_r(0)} \psi_0(r) \frac{\partial \Delta \phi}{\partial r} - \phi(r) \frac{\partial \Delta \psi_0}{\partial r}$$

and hence for r large, we have

$$\frac{\partial \Delta \phi}{\partial r} = -\phi(r) \frac{\partial \Delta \psi_0}{\partial r} = O(\frac{\ln r}{r^5}) \tag{5.11}$$

From (5.9) and (5.11), we see that necessarily, $\int_{\mathbb{R}^4} e^U \phi = 0$. The lemma is thus proved.

Let Π be the stereographic projection of the sphere S^4 onto \mathbb{R}^4 with respect to the North Pole. Namely,

$$\forall x = (x_1, ..., x_5) \in S^4, y_i = \frac{x_i}{1 - x_5}, i = 1, ..., 4, \ \Pi(x) = y = (y_1, ..., y_4)$$
 (5.12)

For a bounded function $\phi(y)$ defined on \mathbb{R}^4 , one can define a function ψ on S^4 by $\psi(x) = \phi(y), y = \Pi(x)$. Then it is easy to see that

$$\int_{S^4} \psi(\sigma) d\sigma = \int_{\mathbb{R}^4} \phi(y) e^U dy \tag{5.13}$$

$$\int_{\mathbb{S}^4} (\mathbb{P}_4 \psi, \psi) = \int_{\mathbb{R}^4} (\Delta \phi)^2 \tag{5.14}$$

where $\mathbb{P}_4 = (-\Delta)(-\Delta + 2)$ is the Paneitz operator on S^4 .

Transforming the identities (5.8) to ψ , we have that ψ satisfies

$$\int_{S^4} \psi d\sigma = \int_{S^4} \psi x_i d\sigma = 0, i = 1, ..., 5$$
(5.15)

Note that the operator $-\Delta_{S^n}$ is known to have eigenvalues λ_j with multiplicity n_j and eigenfunctions u_j as follows:

$$\lambda_0 = 0, n_0 = 1, u_0 = 1,$$

$$\lambda_0 = n, n_1 = n + 1, u_{1,i} = \sigma_i, i = 1, ..., n + 1$$

$$\lambda_k = k(n + k - 1), n_k = \frac{(n + k - 2)!(n + 2k - 1)!}{k!(n - 1)!}, u_{k,i}$$

Thus $\mathbb{P}_4 = (-\Delta)(-\Delta + 2)$ also has eigenvalues and eigenfunctions

$$\mu_0 = 0, u_0 = 1,$$

$$\mu_1 = 24, u_{1,i} = x_i, i = 1, ..., 5$$

$$\mu_2 > 24$$

From (5.6) and (5.14), we derive that

$$\int_{S^4} (\mathbb{P}_4 \psi, \psi) = 24 \int_{\mathbb{R}^4} e^U \phi^2 = 24 \int_{S^4} \psi^2 d\sigma \ge \mu_2 \int_{S^4} \psi^2$$
 (5.16)

which implies $\psi = 0$ and so $\phi = 0$.

This proves Lemma 2.6.

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